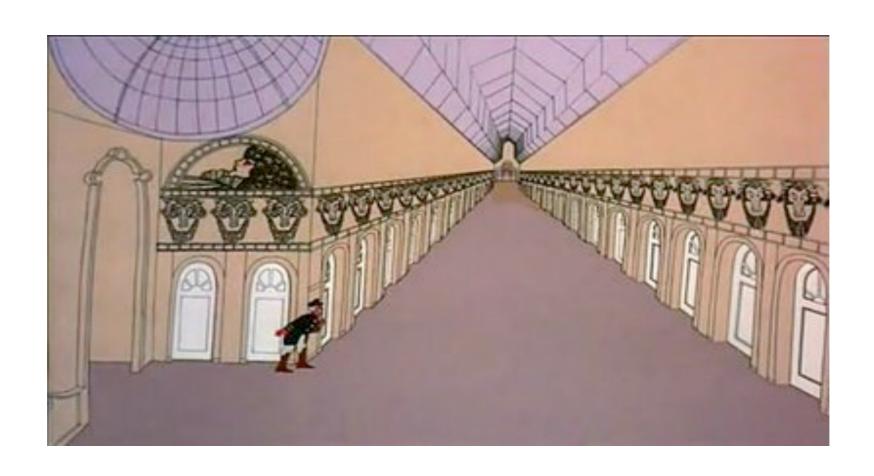
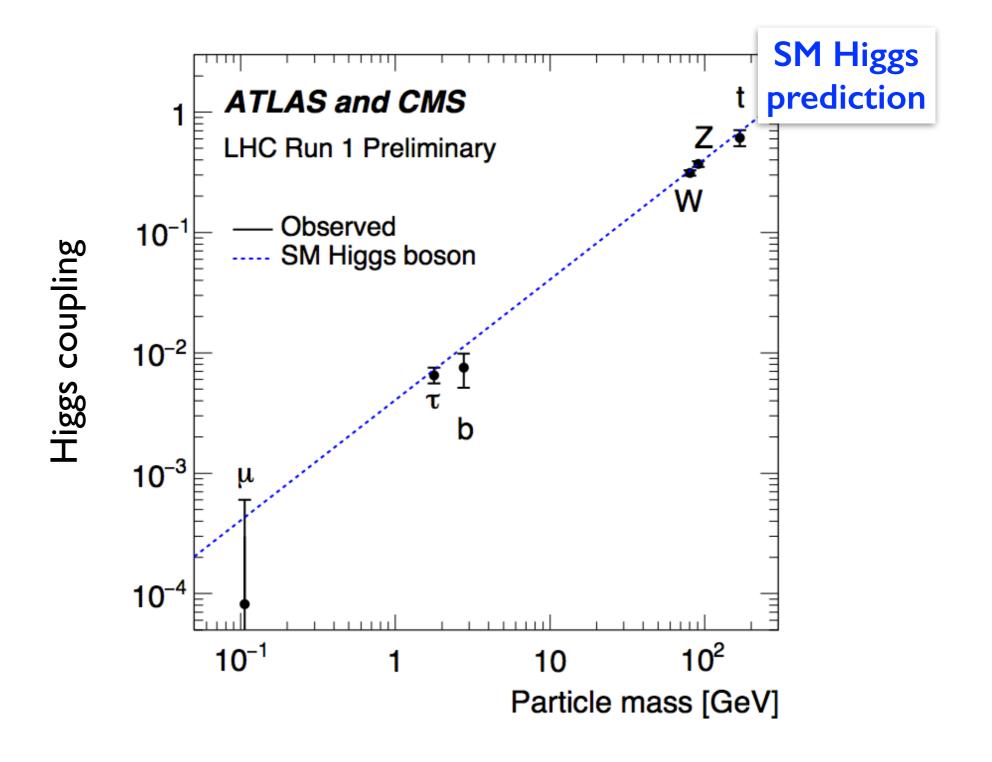
# Searching for new paradigms after the Higgs



Alex Pomarol, CERN & UAB (Barcelona)

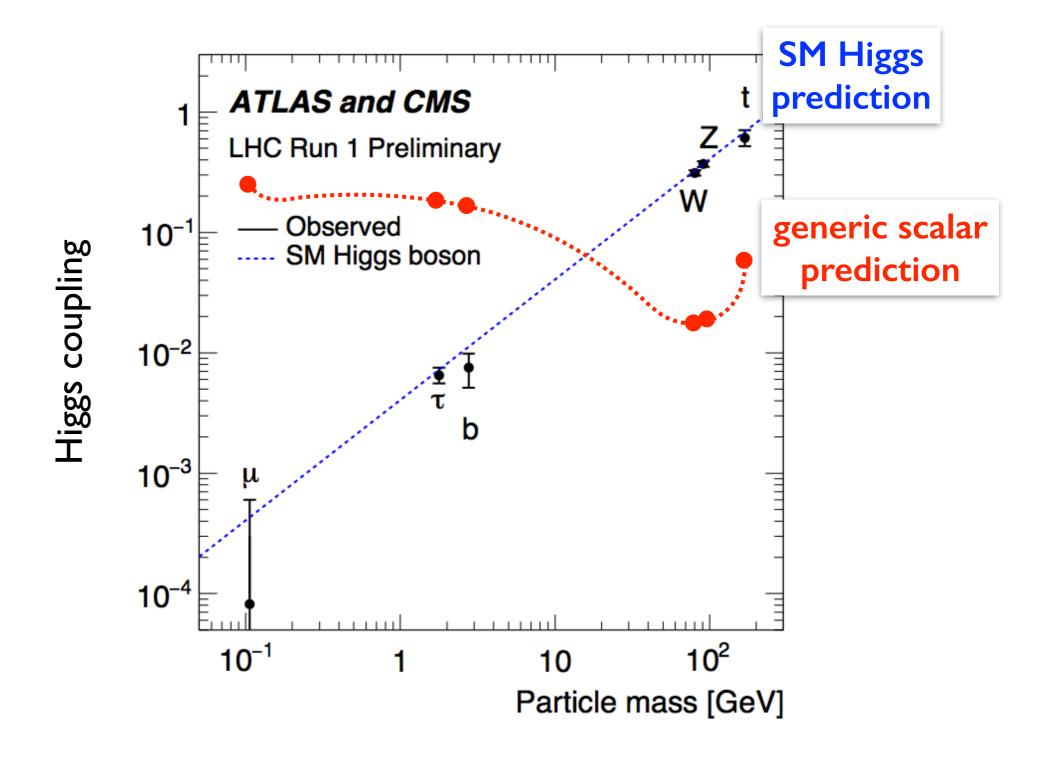
#### Best LHC Run I legacy: The Higgs discovery





Coupling-Mass relations as in the SM Higgs

### Best LHC Run I legacy: The Higgs discovery





"Higgs impostors" left behind!

#### The SM is established!



Where to expect new-physics (beyond the SM)? Where a new paradigm is needed?

#### The SM is established!



# Where to expect new-physics (beyond the SM)? Where a new paradigm is needed?

To answer this, we can follow Einstein's path:

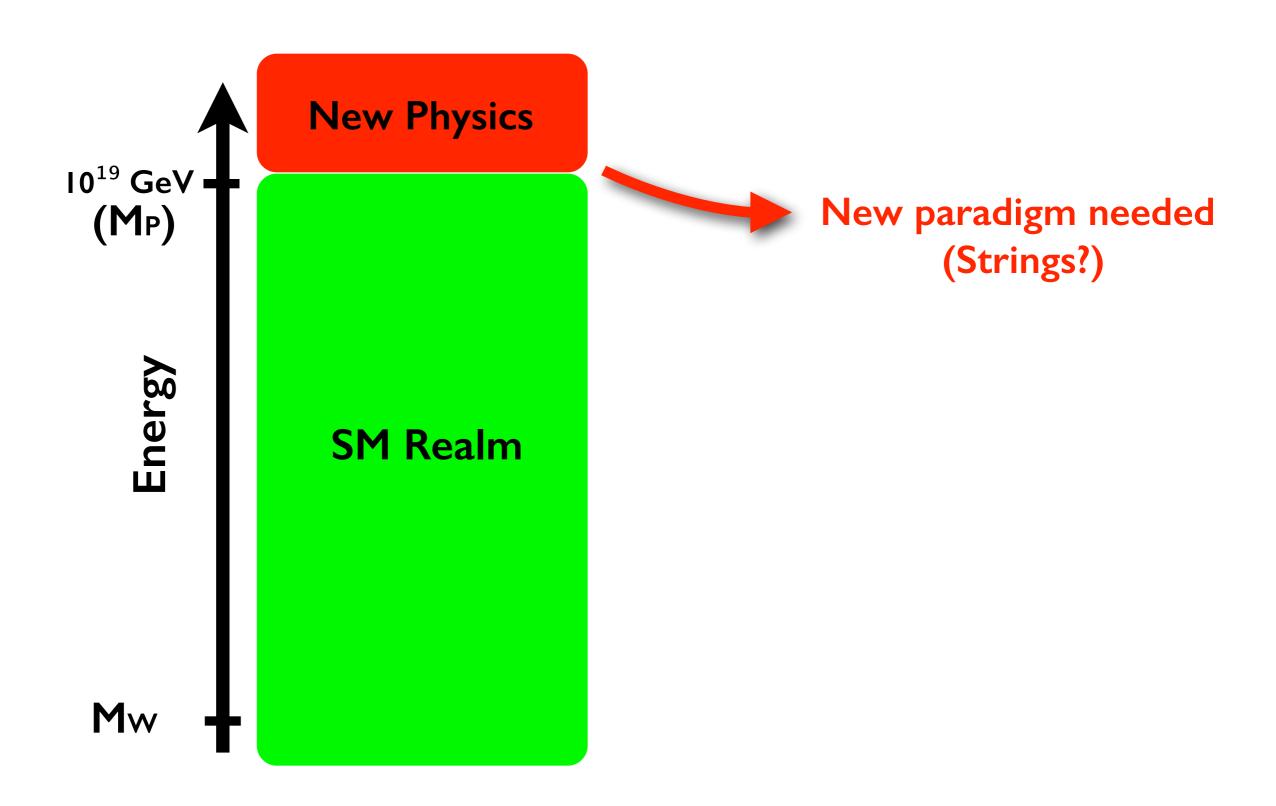


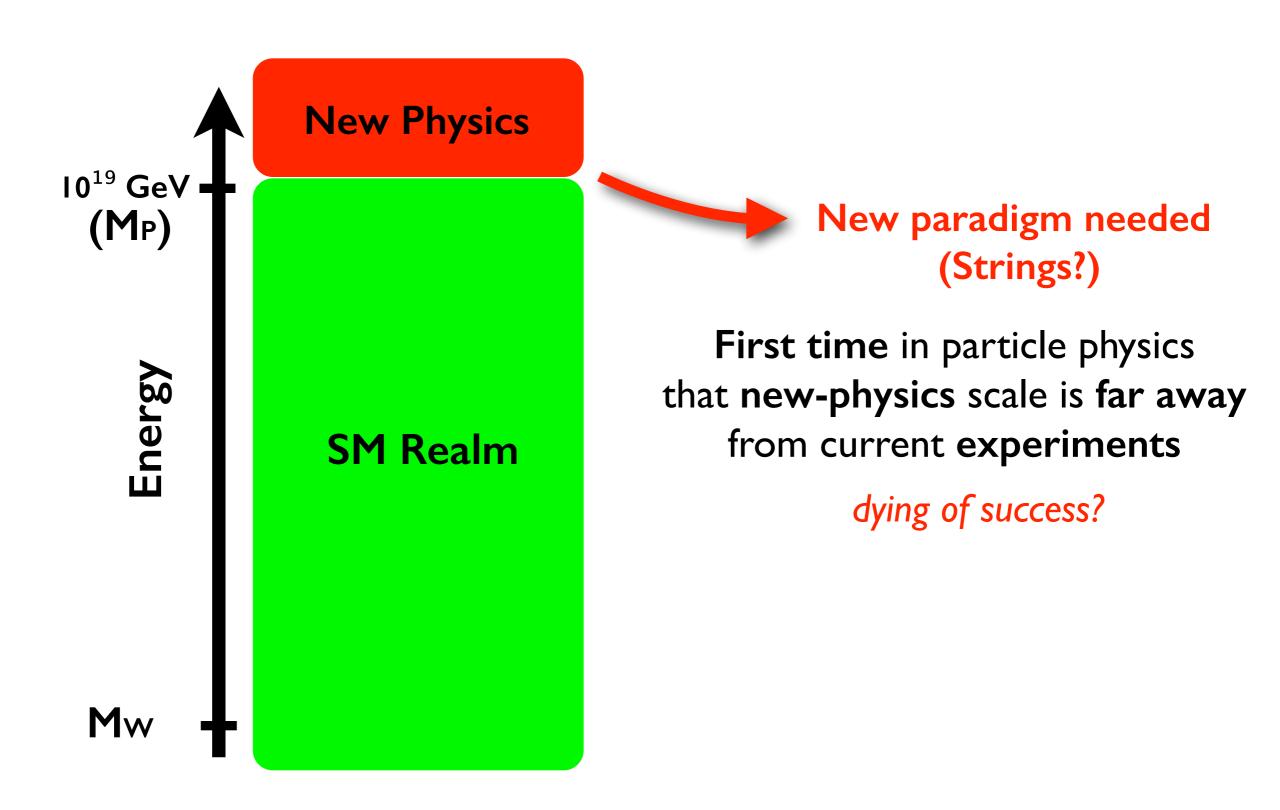
"Gedankenexperiment" (thought experiments):

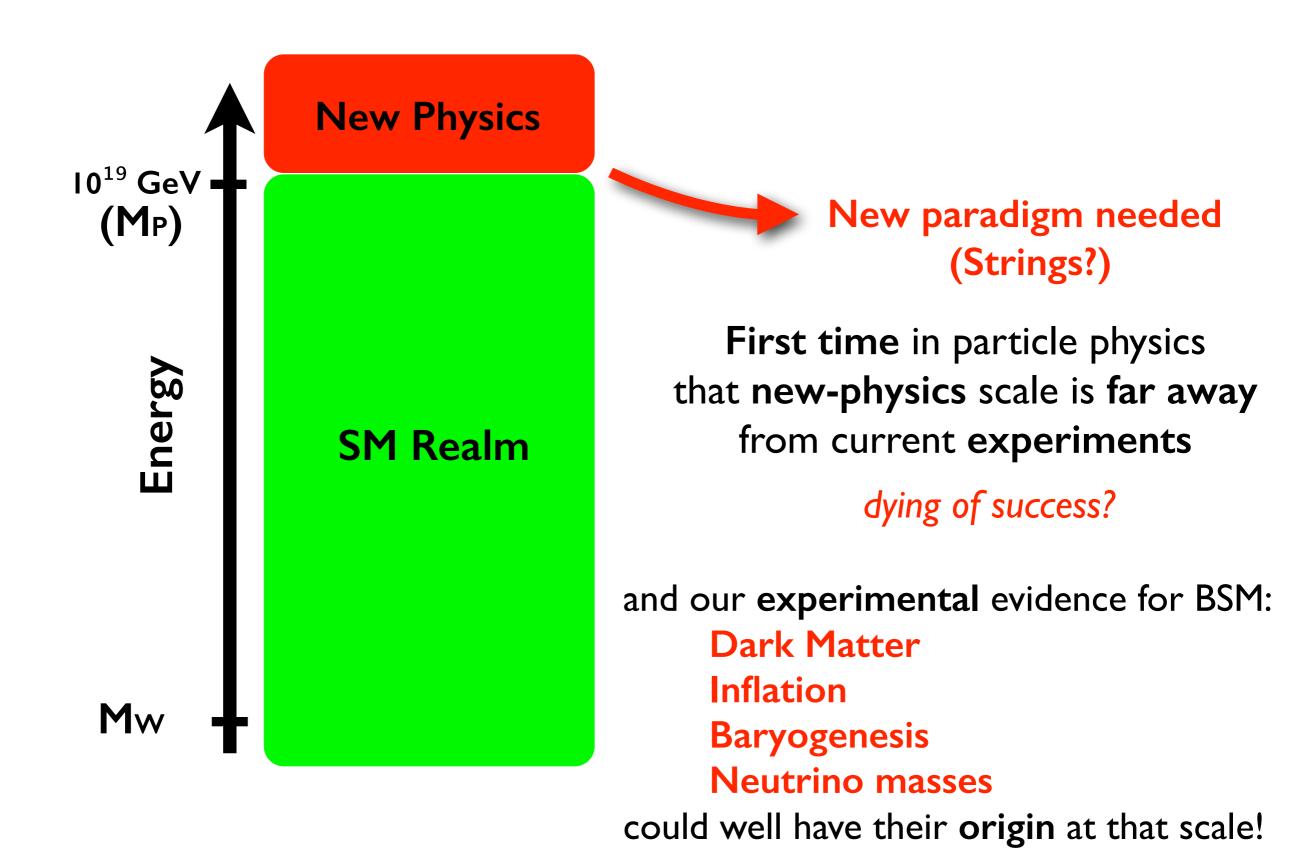
look at which regime the theory fails, and therefore new physics must appear!

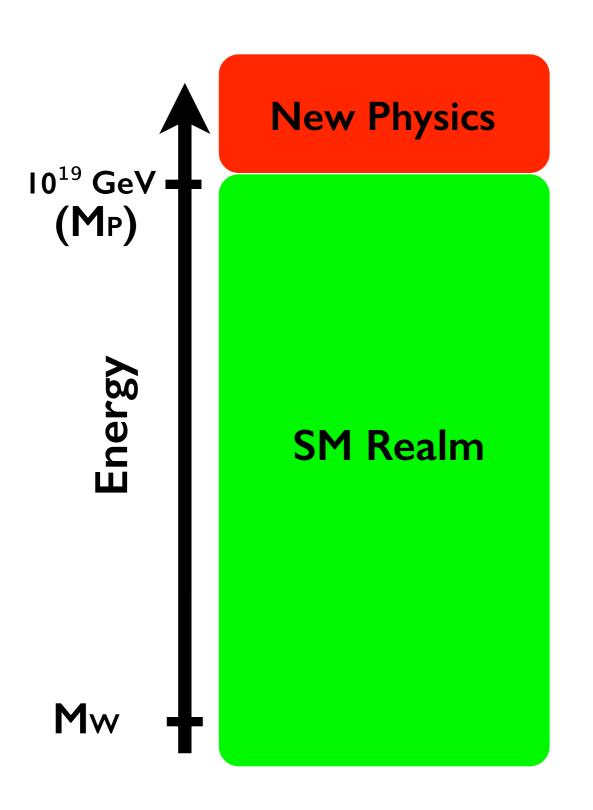
no-lose theorem for a discovery

guaranteed the discover of the positron, charm,..., top & Higgs (or something else)

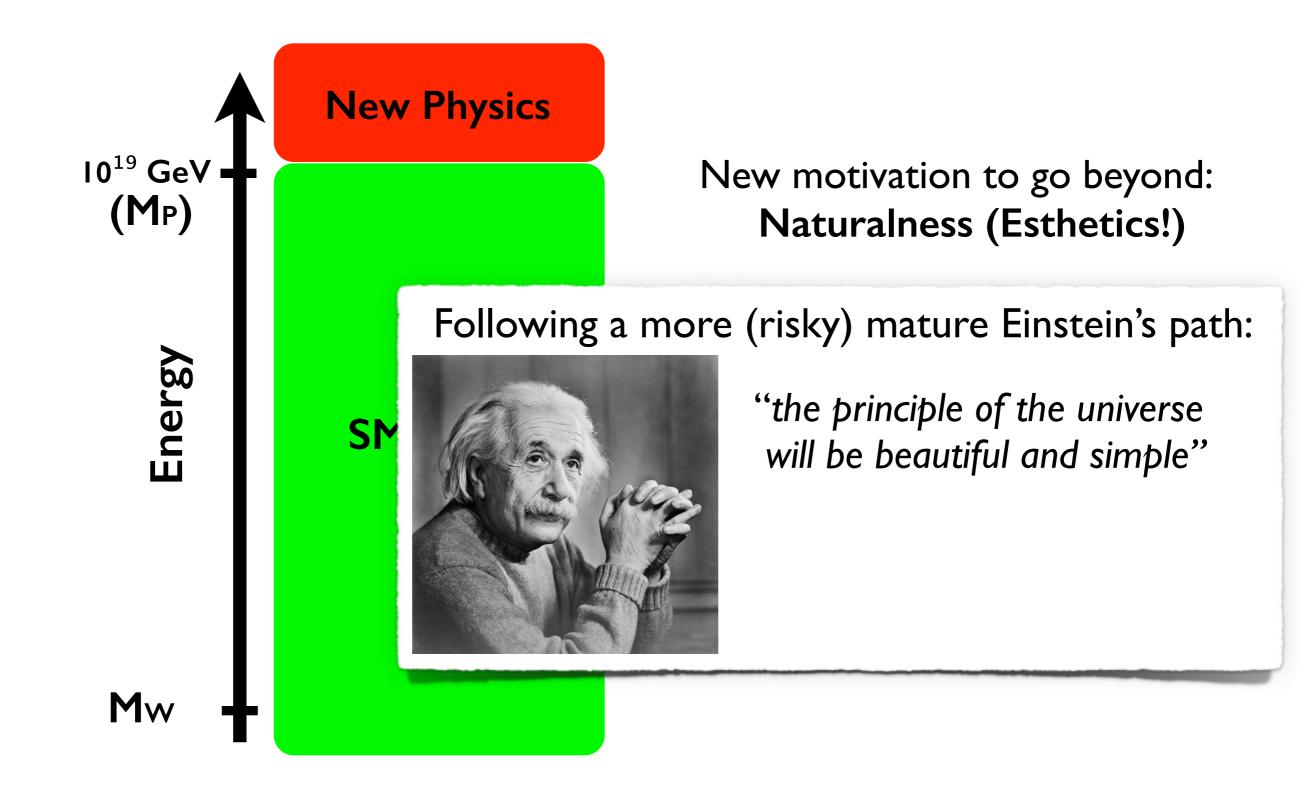


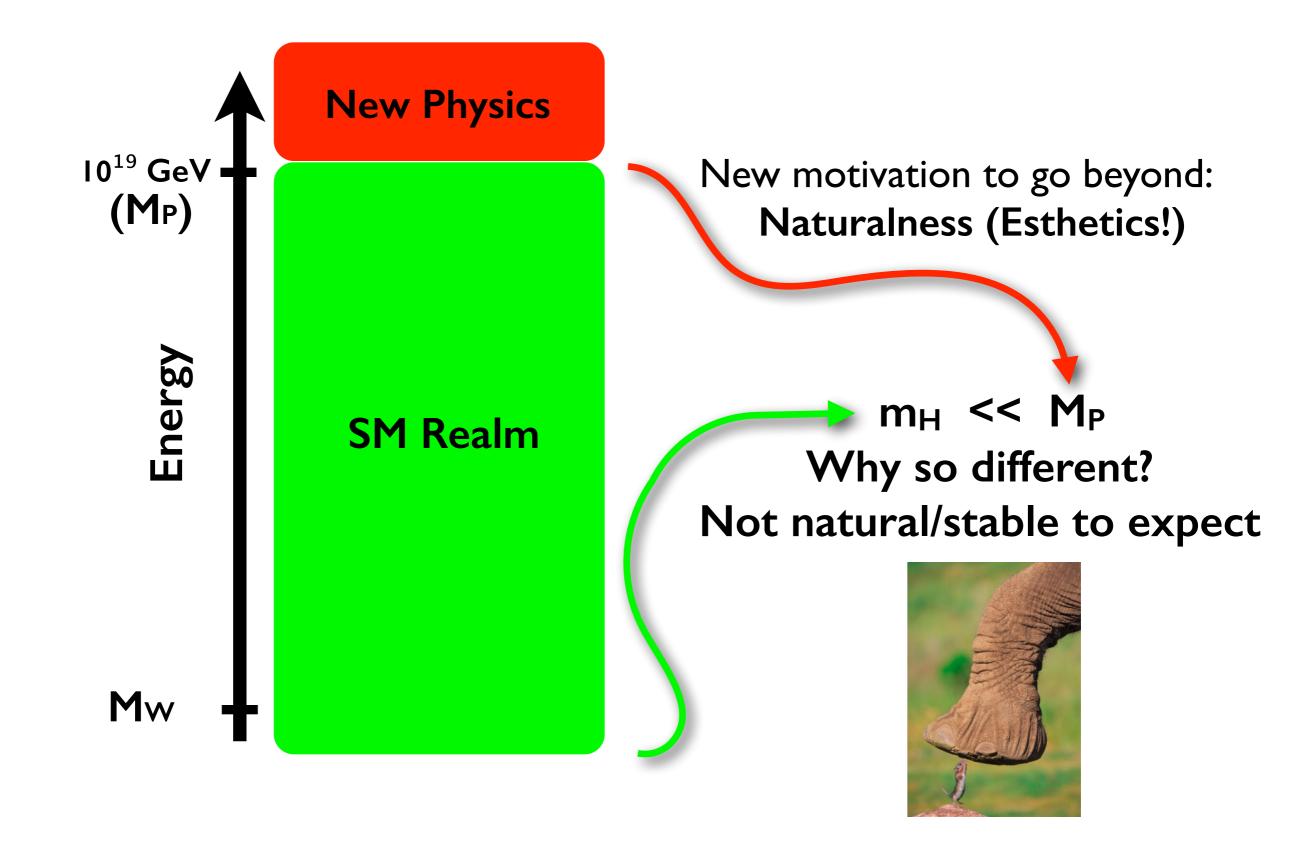






New motivation to go beyond: Naturalness (Esthetics!)





# The Higgs-mass problem in a nutshell

	Massless	Massive	
Vector ${f A}_{\mu}$	2 dof	3 dof	2≠3 ✓ Massless vectors
	(+,-)	(+,0,-)	are save
Fermion	2 dof	4 dof	2≠4 ✓ Massless fermions
Ψ	YL	YL, YR	are save
Scalar h	I dof	I dof	I=I Problem!

### The Higgs-mass problem in a nutshell

Massless Massive

Vector  $\mathbf{A}_{\mu}$ 

2≠3 ✓ Massless vectors are save

Fermion Ψ

2≠4 ✓ Massless fermions are save

Scalar h

I dof

I dof

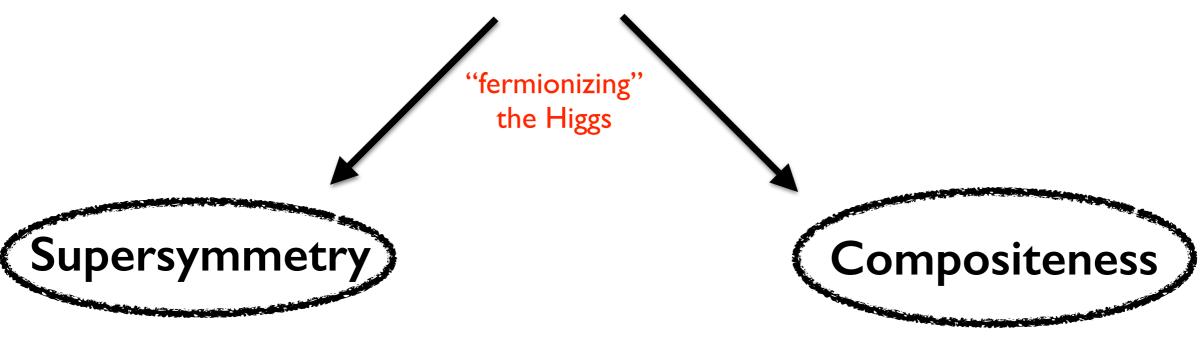
I = I Problem!



Quantum fluctuations can give mass to scalars

# $m_H << M_P$ ?

### Towards a new paradigm



"Marrying" a fermion:

Higss ← → Higgsino

The "transvestite" Higgs:

# Attacking the new paradigm from several fronts

Looking for deviations in **Z/W** couplings

& new particles

LEP/Tevatron









Looking for deviations in **Higgs** couplings & new particles





Looking for Electric Dipole Moments

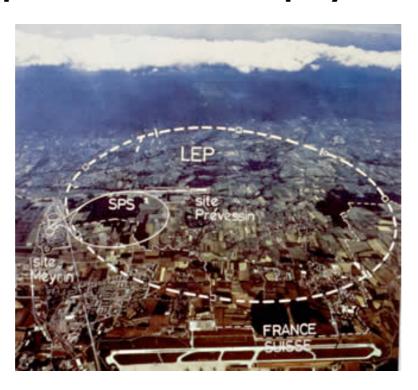


No Success so far!

new-physics

# First main weapon to attack physics Beyond the SM (BSM):

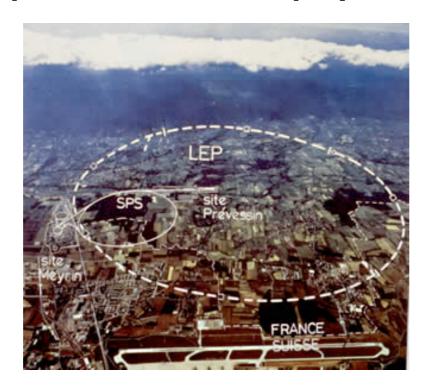




~ millions of Z produced

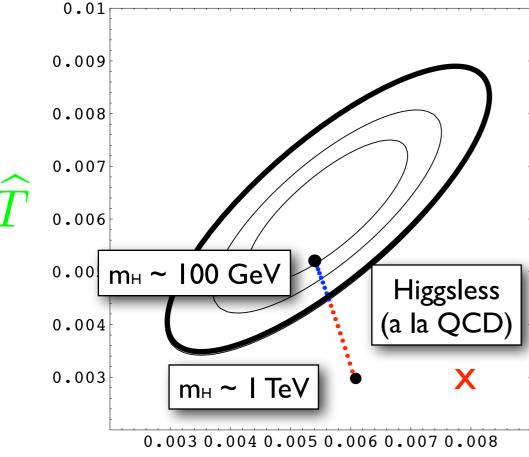
#### First main weapon to attack physics Beyond the SM (BSM):





~ millions of Z produced





Deviations were expected at the >1% level
But no sign of New-physics!



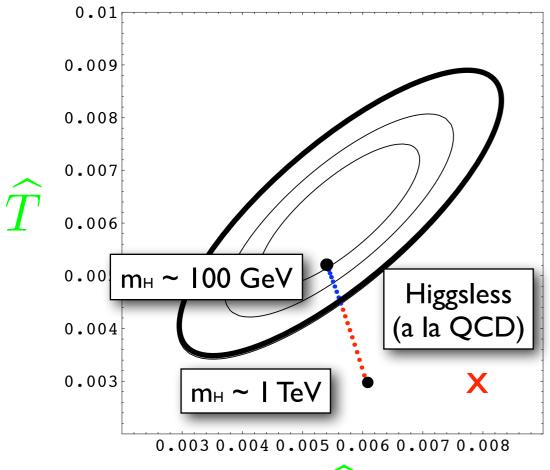
#### First main weapon to attack physics Beyond the SM (BSM):





~ millions of Z produced



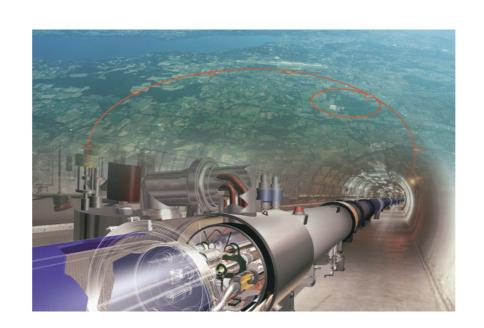


Deviations were expected at the >1% level
But no sign of New-physics!

**Bad luck?** 

We built a more powerful weapon:

# LHC



It has brought an important new discovery: The Higgs!



With the **Higgs**, we have had access to new relevant information by measuring its **properties** 

The Higgs is the most "sensitive"

SM particle to new-physics,

and therefore

the best place to look for natural BSM

#### **Examples:**

Superpartners

Z

Gauge bosons:

A ree-level effects

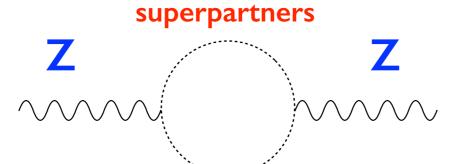
#### **Examples:**

#### **Examples:**

#### I) MSSM:

Gauge bosons:

Higgs:



~ loop effects

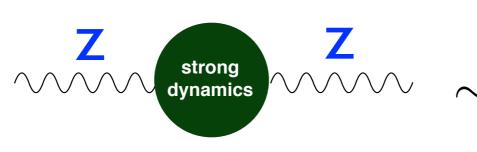
h H ~ tree-level effects



Effects in Higgs physics can be a factor  $16\pi^2 \sim 100$  larger!

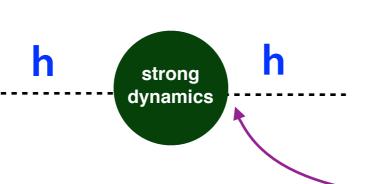
#### 2) Composite models:

Gauge bosons:



 $(\Lambda = composite scale)$ 

Higgs:



 $\frac{g_H^2 v^2}{\Lambda^2} \sim \frac{16\pi^2 v^2}{\Lambda^2}$ 

-"strong" Higgs coupling

#### Consequences:

Even with less statistics at the LHC, similar impact today in new-physics as LEP

LEP:  $ee \rightarrow Z (\rightarrow ff) \sim millions of events$ 

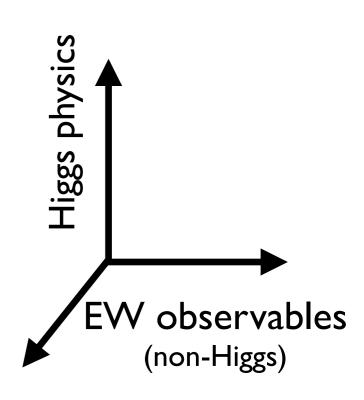
LHC:  $pp \rightarrow h (\rightarrow \gamma \gamma) \sim thousands of events$ 

#### First question to address in Higgs couplings:

# Which are the most relevant Higgs couplings to measure?

probes testing
new directions in the
"parameter space" of BSMs

Couplings that can be modified by new-physics, not affecting anything else



#### Assuming new-physics scale is heavier than m<sub>H</sub>:

# 8 Primary Higgs couplings

(for CP-conservation and one family)

Elias-Miro, Espinosa, Masso, AP, JHEP 1311 (2013) 066 AP, Riva, JHEP 1401 (2014) 151

$$\Delta \mathcal{L}_{\text{BSM}} = \frac{\delta g_{hff}}{h \bar{f}_L f_R + h.c.} \qquad \text{(f=b, \tau, t)}$$

$$+ \frac{g_{hVV}}{g_{hVV}} h \left[ W^{+\mu} W_{\mu}^{-} + \frac{1}{2 \cos^2 \theta_W} Z^{\mu} Z_{\mu} \right]$$

$$+ \frac{\kappa_{GG}}{v} G^{\mu\nu} G_{\mu\nu}$$

$$+ \frac{\kappa_{\gamma\gamma}}{v} \frac{h}{v} F^{\gamma \mu\nu} F_{\mu\nu}^{\gamma}$$

$$+ \frac{\kappa_{\gamma Z}}{v} \frac{h}{v} F^{\gamma \mu\nu} F_{\mu\nu}^{Z}$$

$$+ \frac{\delta g_{3h}}{v} h^3$$

Assuming new-physics scale is heavi

# 8 Primary Higgs couplings

$$\Delta \mathcal{L}_{\text{BSM}} = \frac{\delta g_{hff}}{\delta g_{hff}} h \bar{f}_L f_R + h.c.$$

$$+ \frac{g_{hVV}}{\delta g_{hV}} h \left[ W^{+\mu} W_{\mu}^{-\mu} \right] + \frac{\kappa_{GG}}{v} G^{\mu\nu} G_{\mu\nu}$$

$$+ \frac{\kappa_{\gamma\gamma}}{v} \frac{h}{v} F^{\gamma\mu\nu} F_{\mu\nu}^{\gamma}$$

$$+ \frac{\kappa_{\gamma Z}}{v} \frac{h}{v} F^{\gamma\mu\nu} F_{\mu\nu}^{Z}$$

$$+ \frac{\delta g_{3h}}{v} h^3$$

Corresponds to the 8 possible dim-6 operators with |H|<sup>2</sup>:

 $|H|^2 \bar{f}_L H f_R + h.c.$  $|H|^2|D_\mu H|^2$  $|H|^2 G^A_{\mu\nu} G^{A\mu\nu}$  $|H|^2 B_{\mu\nu} B^{\mu\nu}$  $|H|^2 W^a_{\mu\nu} W^{\mu\nu\,a}$  $|H|^{6}$ 

66



#### Assuming new-physics scale is heavier than m<sub>H</sub>:

# 8 Primary Higgs couplings

(for CP-conservation and one family)

Elias-Miro, Espinosa, Masso, AP, JHEP 1311 (2013) 066 AP, Riva, JHEP 1401 (2014) 151

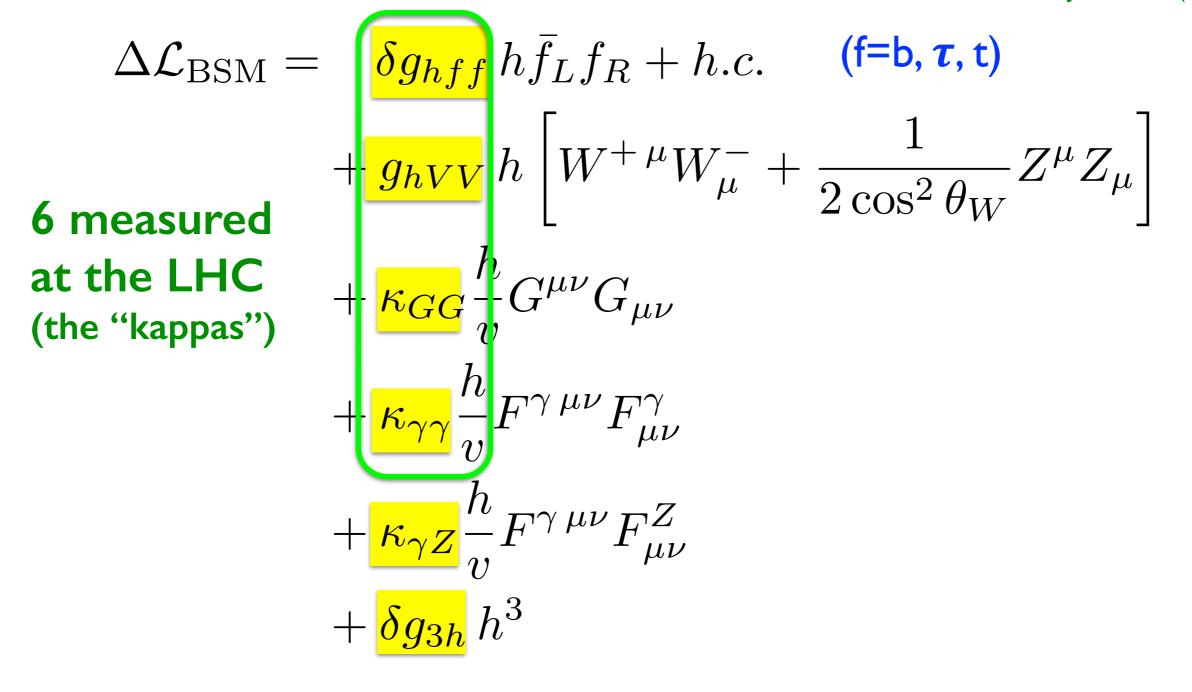
$$\begin{split} \Delta \mathcal{L}_{\mathrm{BSM}} = & & \delta g_{hff} \, h \bar{f}_L f_R + h.c. \quad \text{(f=b, $\tau$, t)} \\ & + \frac{g_{hVV}}{g_{hVV}} \, h \left[ W^{+\,\mu} W_\mu^- + \frac{1}{2 \cos^2 \theta_W} Z^\mu Z_\mu \right] \\ & + \frac{\kappa_{GG}}{v} G^{\mu\nu} G_{\mu\nu} \\ & + \frac{\kappa_{\gamma\gamma}}{v} \frac{h}{v} F^{\gamma\,\mu\nu} F^{\gamma}_{\mu\nu} & \text{important: custodial invariant!!} \\ & + \frac{h}{\kappa_{\gamma Z}} \frac{h}{v} F^{\gamma\,\mu\nu} F^Z_{\mu\nu} \\ & + \frac{\delta g_{3h}}{v} \, h^3 \end{split}$$

#### Assuming new-physics scale is heavier than m<sub>H</sub>:

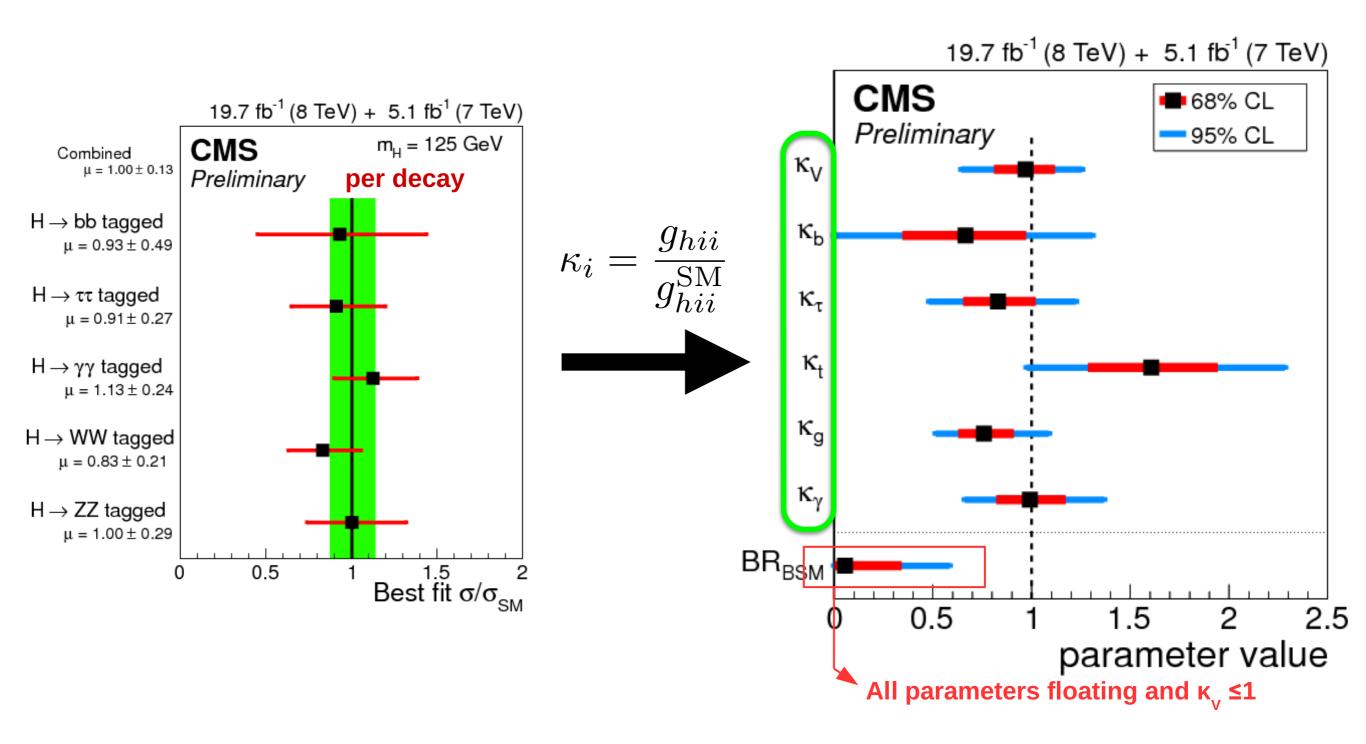
# 8 Primary Higgs couplings

(for CP-conservation and one family)

Elias-Miro, Espinosa, Masso, AP, JHEP 1311 (2013) 066 AP, Riva, JHEP 1401 (2014) 151

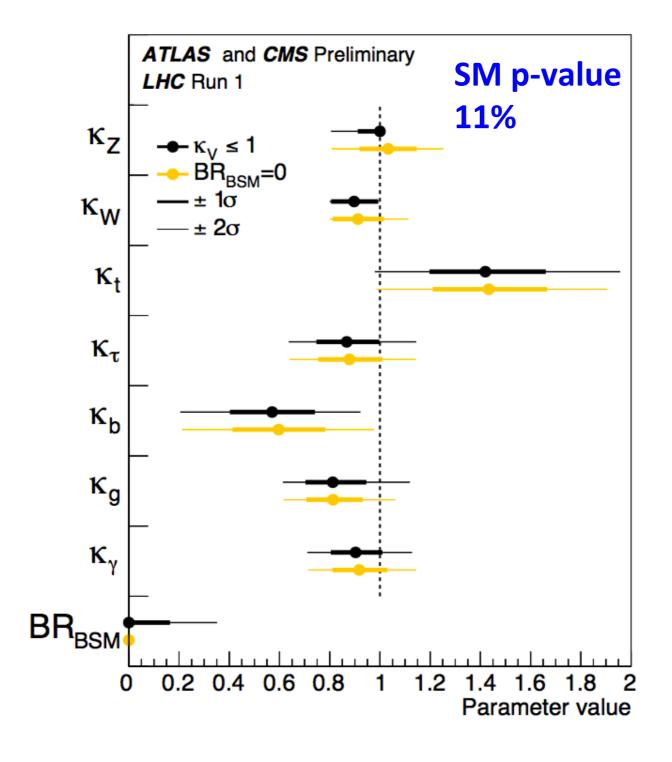


# Higgs coupling determination



reasonable good agreement with the SM!

#### Combined analysis:



reasonable good agreement with the SM!

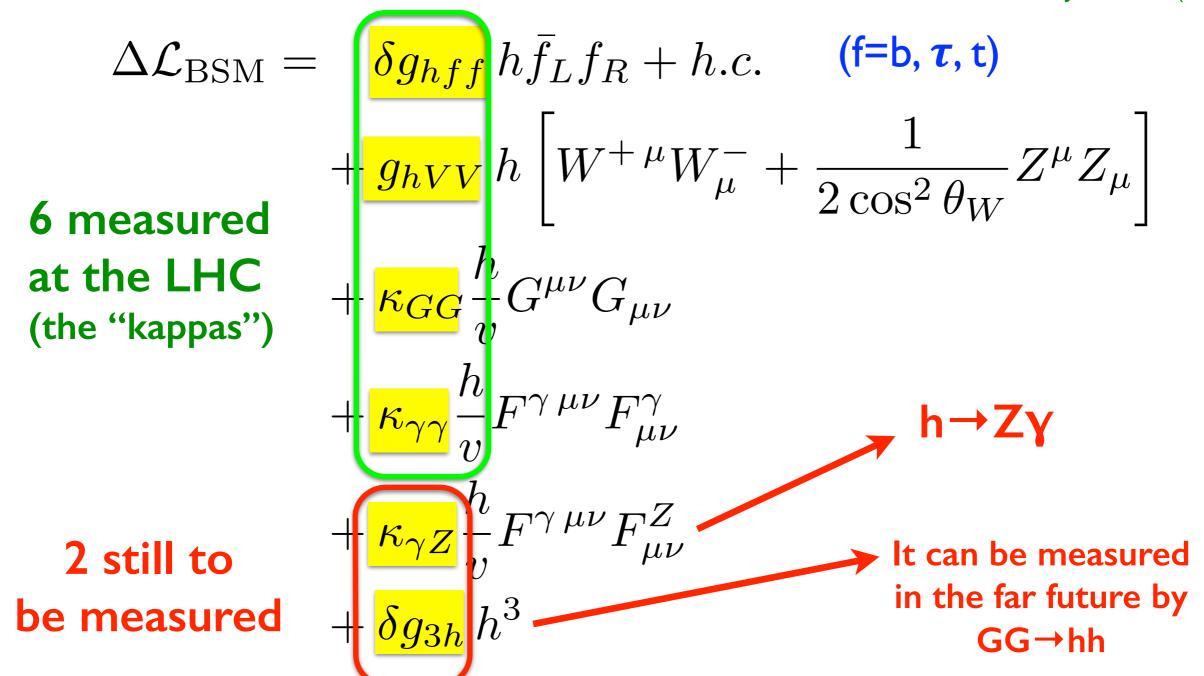


#### Assuming new-physics scale is heavier than m<sub>H</sub>:

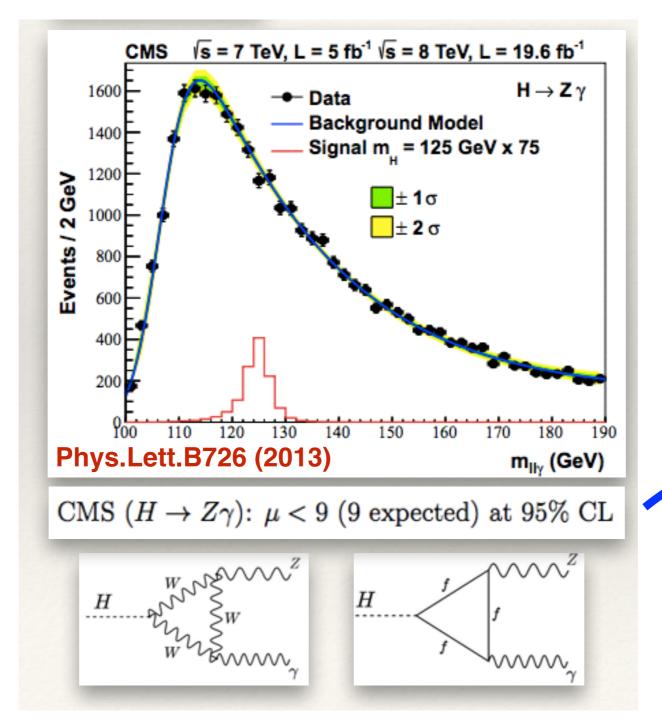
# 8 Primary Higgs couplings

(for CP-conservation and one family)

Elias-Miro, Espinosa, Masso, AP, JHEP 1311 (2013) 066 AP, Riva, JHEP 1401 (2014) 151



# Experimental bound on $h \rightarrow Z\gamma$



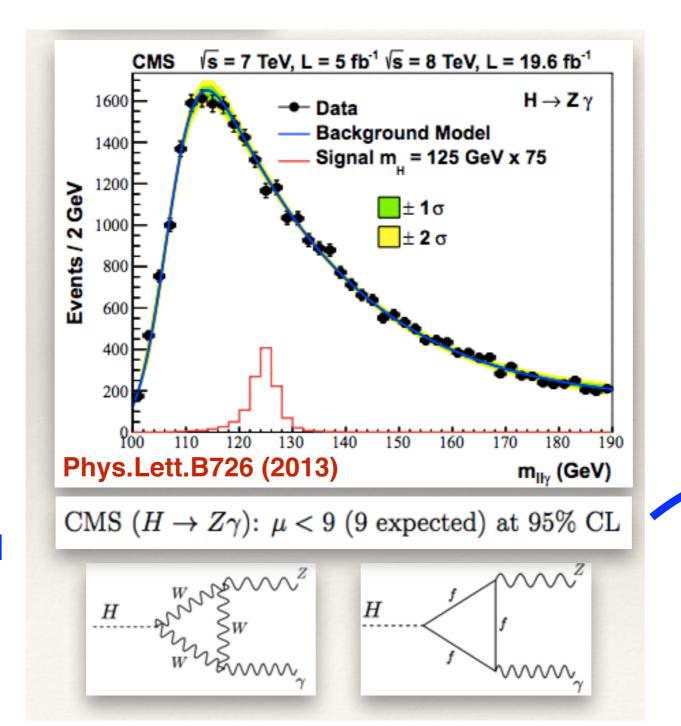
still allowed to be 9 x BR<sub>SM</sub>

BR( $h\rightarrow Z\gamma$ )~0.001 small in the SM since it comes at one-loop:

#### ... last hope for finding O(I) deviations?

(possibility in composite Higgs models)

# Experimental bound on $h \rightarrow Z\gamma$



BR( $h\rightarrow Z\gamma$ )~0.001 small in the SM since it comes at one-loop:

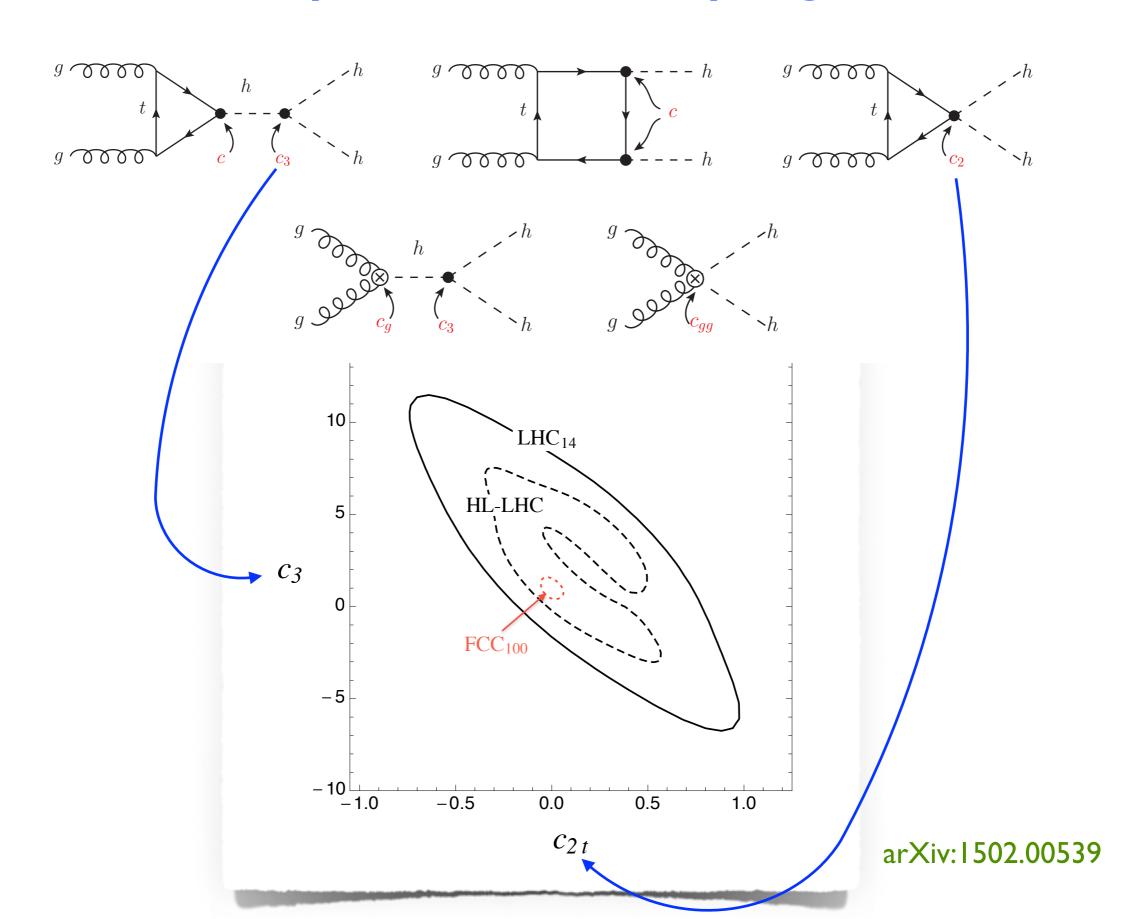
still allowed to be 9 x BR<sub>SM</sub>

... last hope for finding O(1) deviations?

(possibility in composite Higgs models)

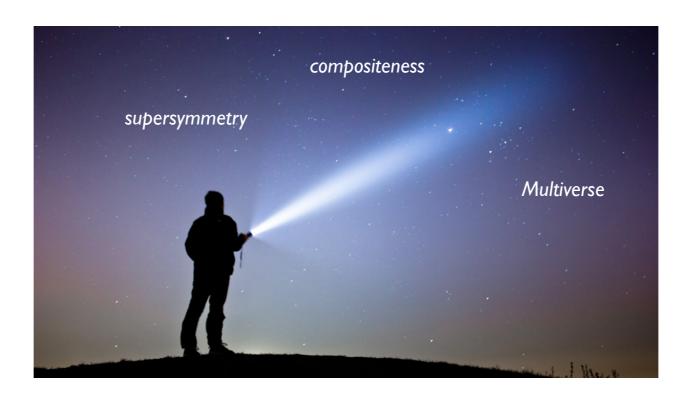


# **Prospects for 3h-coupling**



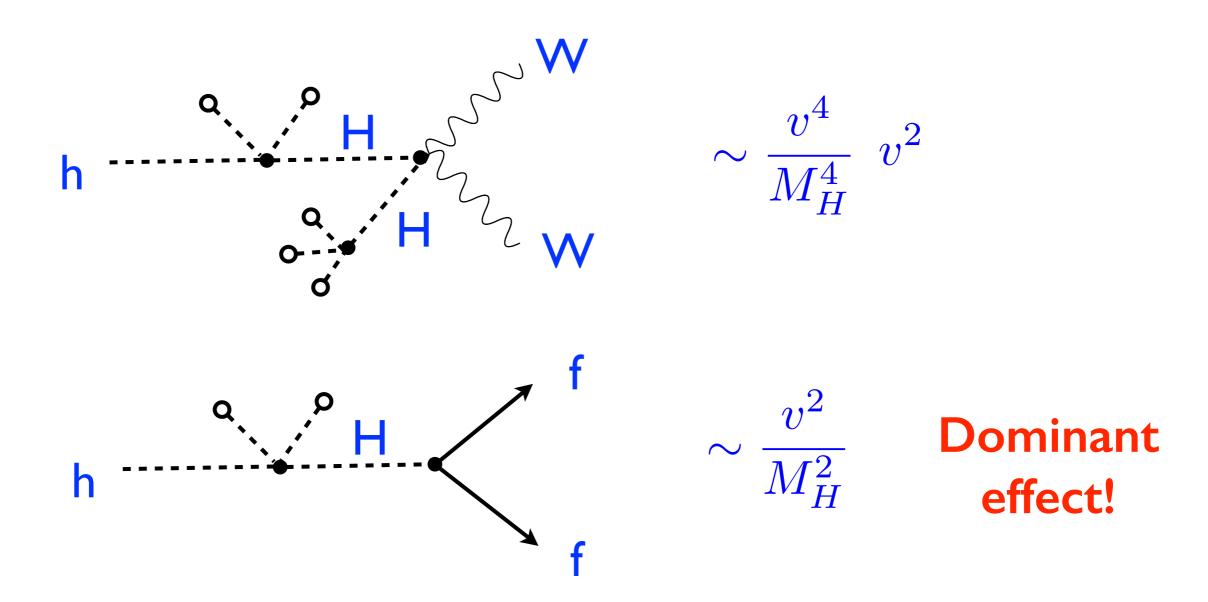
# Impact on BSM from Higgs coupling measurements

- Today, as Higgs coupling measurements agree with the SM, we only place bounds on new-physics
  - The Higgs is our best weapon of BSM mass-destruction
- Tomorrow, who knows, it can illuminate on new-physics



# MSSM with heavy spectrum ( > 100 GeV)

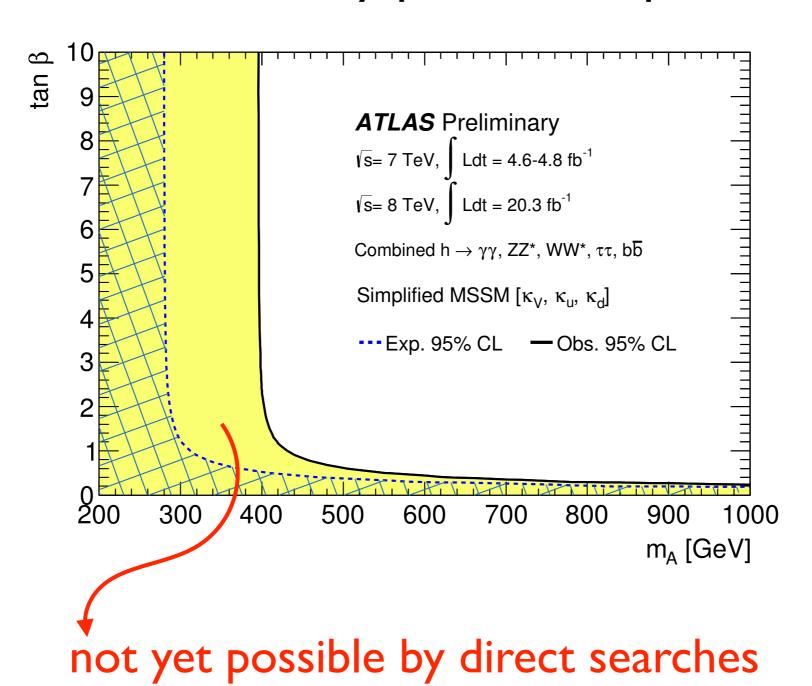
#### Main effects from the 2nd Higgs doublet:



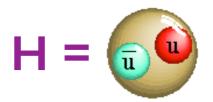
Superpartners can only modify Higgs couplings at the loop-level: Only stops/sbottoms give some contribution to hgg/h $\gamma\gamma$  (not very large)

## Supersymmetric Models (MSSM)

Higgs coupling measurements already rules out susy-parameter space



## **Composite Higgs**



(Higgs as a pion)

Couplings dictated by symmetries (as in the QCD chiral Lagrangian)

$$\frac{g_{hWW}}{g_{hWW}^{SM}} = \sqrt{1 - \frac{v^2}{f^2}}$$

Giudice, Grojean, AP, Rattazzi 07

f = Decay-constant of the PGB Higgs related to the compositeness scale
(model dependent but expected f ~ v)

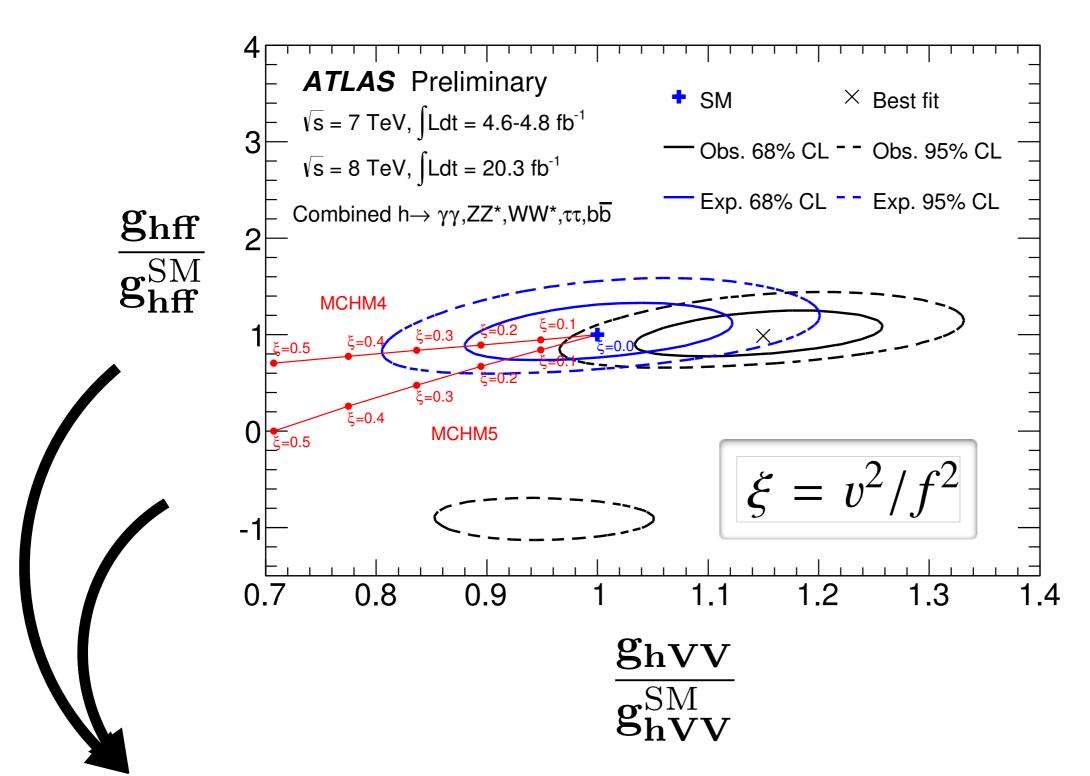
$$\frac{g_{hff}}{g_{hff}^{SM}} = \frac{1 - (1+n)\frac{v^2}{f^2}}{\sqrt{1 - \frac{v^2}{f^2}}}$$

AP,Riva 12

$$n = 0, 1, 2, \dots$$

$$\uparrow \qquad \uparrow$$
MCHM4 MCHM5

small deviations on the  $h\gamma\gamma$ (gg)-coupling due to the Goldstone nature of the Higgs



observed (expected) 95% CL upper limit of  $\xi$  < 0.12 (0.29) **MCHM4**  $\xi$  < 0.15 (0.20) **MCHM5** 

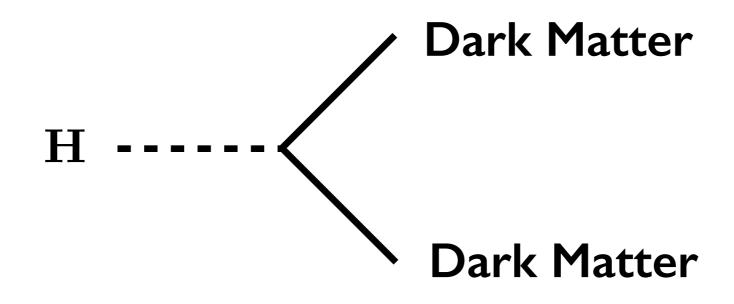
The most interesting one:

I) invisible Higgs decay:



## The most interesting one:

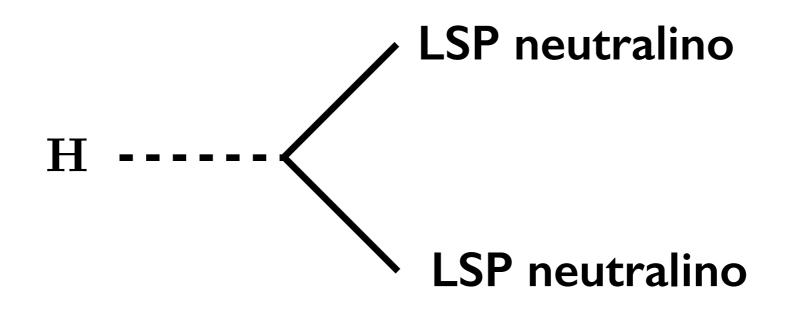
I) invisible Higgs decay:





## The most interesting one:

I) invisible Higgs decay:

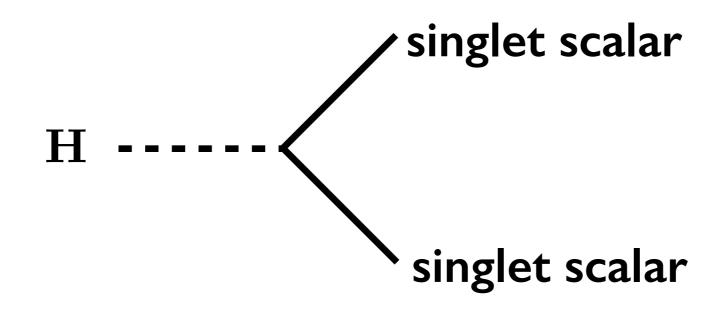




in supersymmetric models

## The most interesting one:

I) invisible Higgs decay:

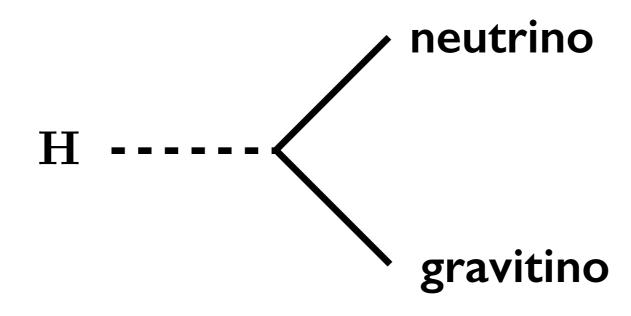




in models with more composite Higgs

## The most interesting one:

I) invisible Higgs decay:

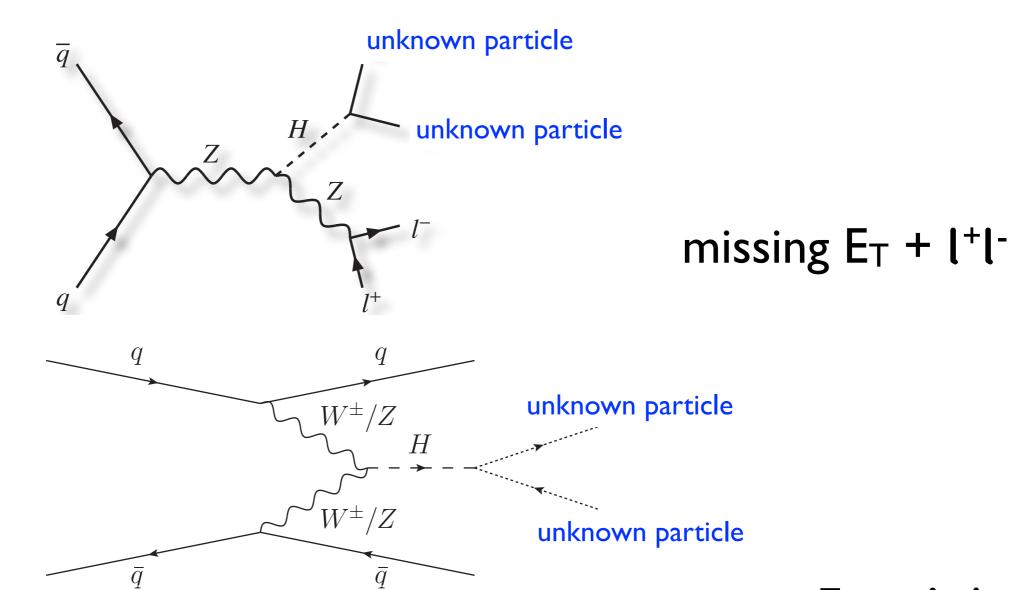




in theories where the Higgs is the superpartner of the neutrino Fayet,'76; AP,Riva,Biggio'12

#### How to "see" it?

**HV** channel:



**VBF** channel:

missing  $E_T$  + jets

No sign of so, up to now:

CMS:  $BR_{inv} < 58\%$  (44% expected)

ATLAS:  $BR_{inv} < 29\%$  (35% expected)

# The LHC is a discovery machine, and the search for new particles is the best option for discovering new physics beyond the SM

(Higgs coupling measurements can only be complementary)

# The LHC is a discovery machine, and the search for new particles is the best option for discovering new physics beyond the SM

(Higgs coupling measurements can only be complementary)

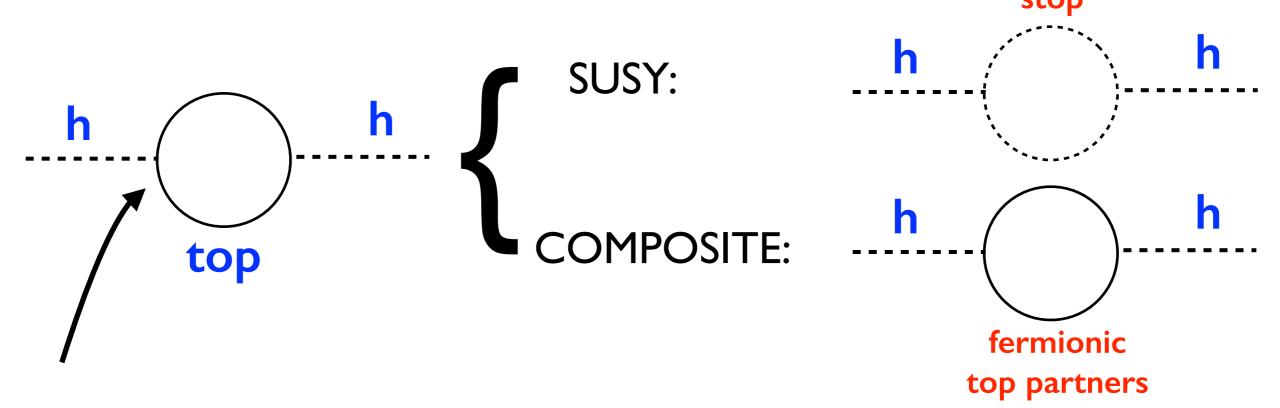
The Higgs tells us where to look for, since new particles must be there to stabilize the Higgs mass!

# The LHC is a discovery machine, and the search for new particles is the best option for discovering new physics beyond the SM

(Higgs coupling measurements can only be complementary)

The Higgs tells us where to look for,

since new particles must be there to stabilize the Higgs mass!



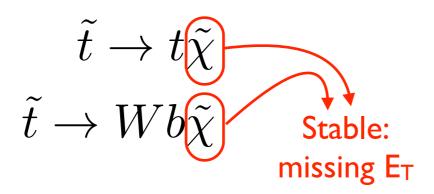
largest Higgs coupling, largest quantum destabilization

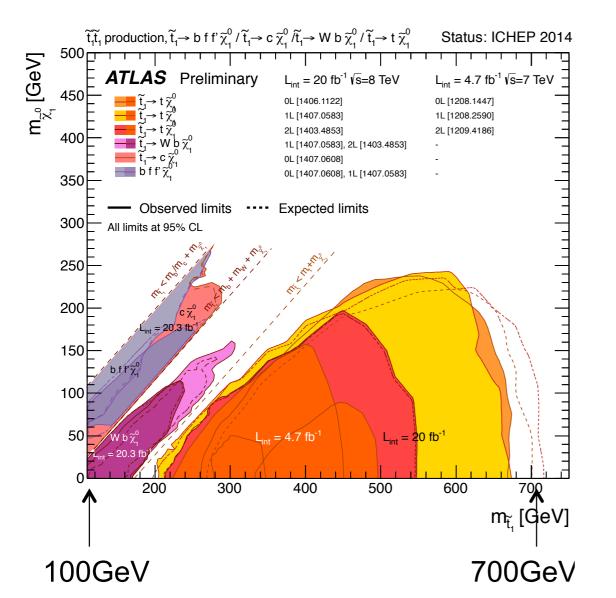
Extra particles to stabilize it

They must be the lightest if no tunings

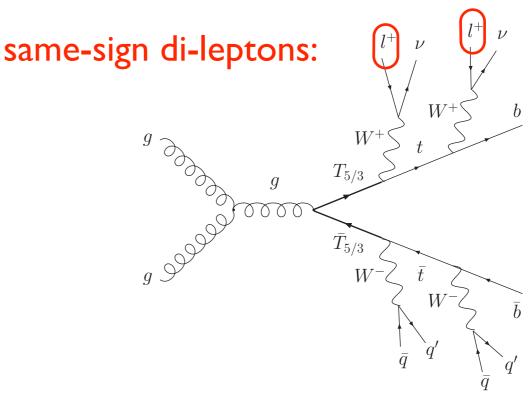
#### Best chances at the LHC to find new physics:

#### **Stops/Sbottoms:**





# Color vector-like fermions with charge 5/3:



ATLAS-CONF-2012-130:

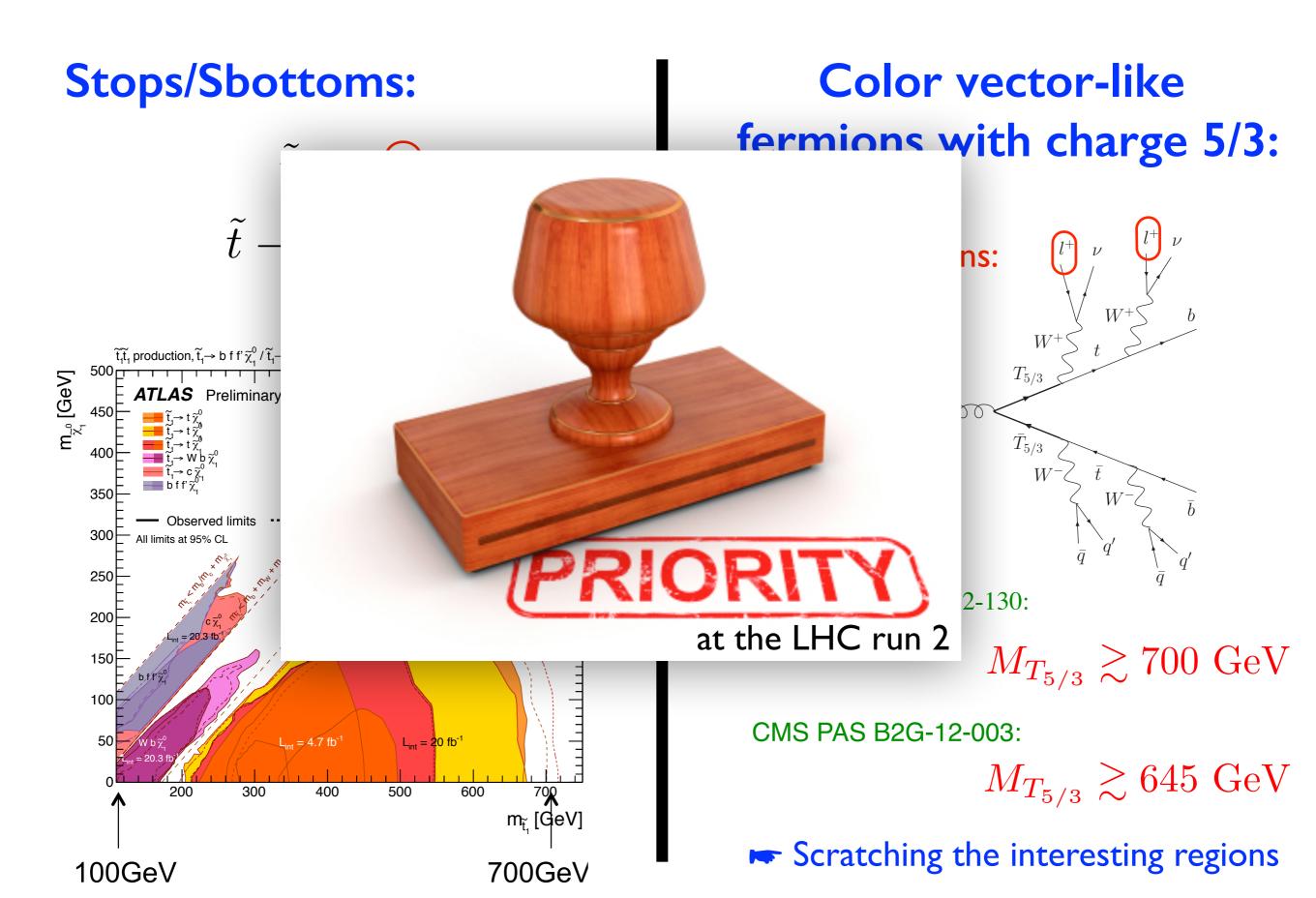
$$M_{T_{5/3}} \gtrsim 700 \text{ GeV}$$

CMS PAS B2G-12-003:

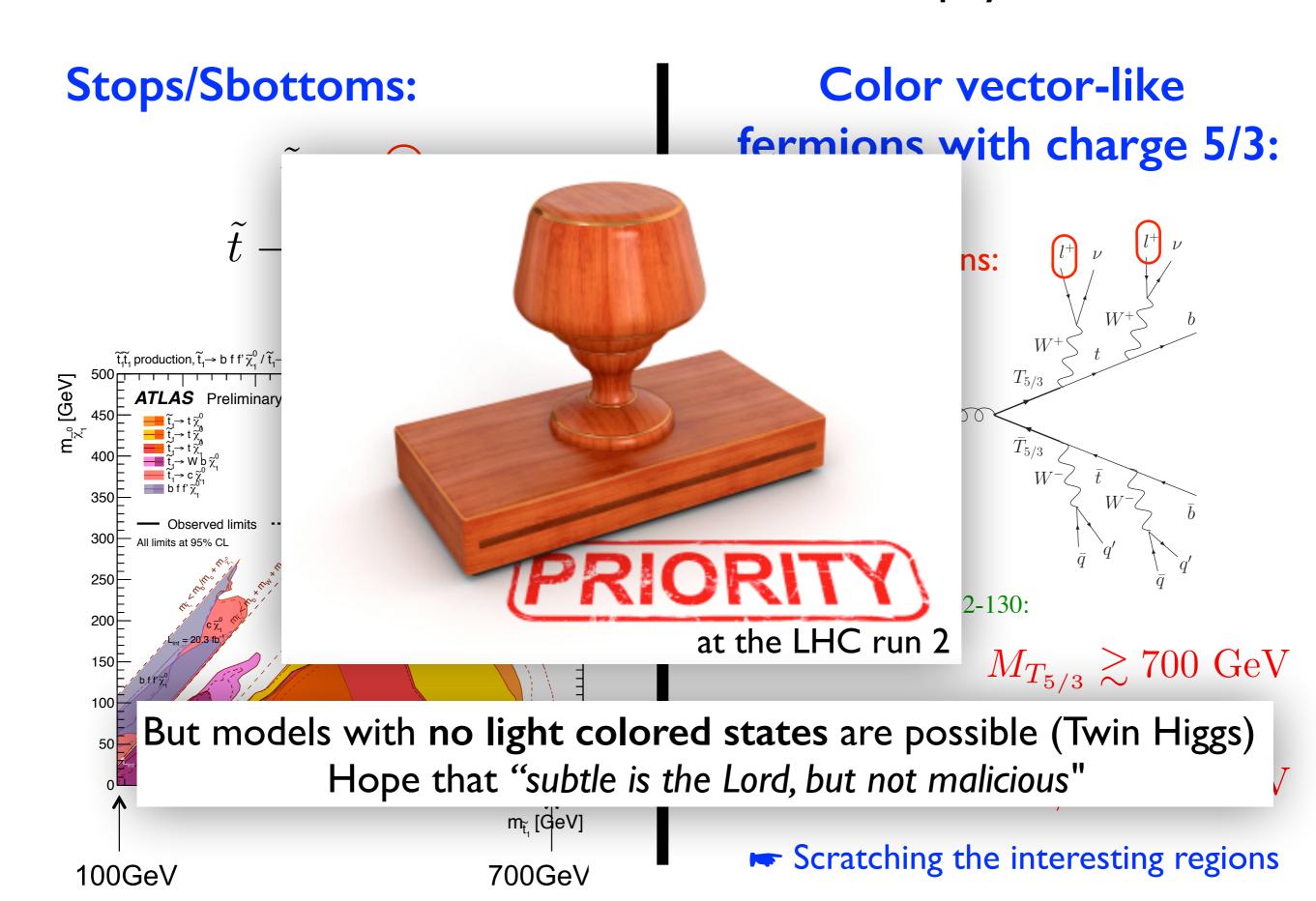
$$M_{T_{5/3}} \gtrsim 645 \text{ GeV}$$

Scratching the interesting regions

### Best chances at the LHC to find new physics:



#### Best chances at the LHC to find new physics:



**MANNY MAN** 

Acceptance imes Efficiency

0.2

#### Di-Boson Resonances

• Focusing on VV production with V=W ( $\rightarrow$  ev, $\mu\nu$ ,qq'), Z( $\rightarrow$  ee, $\mu\mu$ ,qq) resonances, skipping possible Higgs combinations:

Sequential Standard Model (SSM), Extended Gauge Model (EGM) W' → WZ decays,

• RS graviton  $G \rightarrow WW$ ,  $G \rightarrow ZZ$  decay.

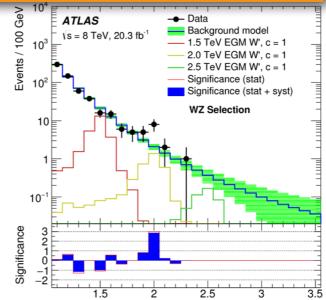
W/Z

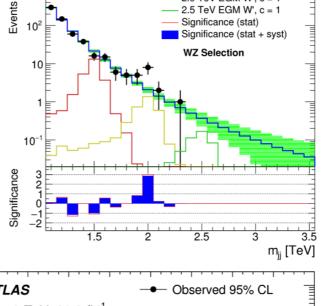
Again, boosted topologies for V→qq

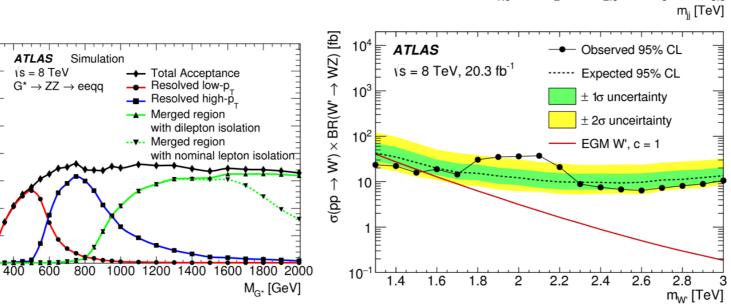
at high mass resonances

... an (3.4  $\sigma$  local, 2.5  $\sigma$  global) excess observed at M ~ 2 TeV by ATLAS in the

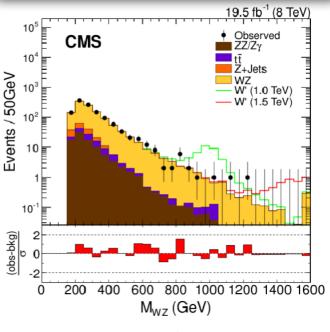
 $VV \rightarrow qq qq final state (only)$ 

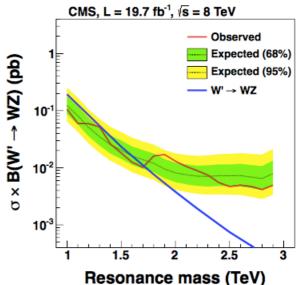




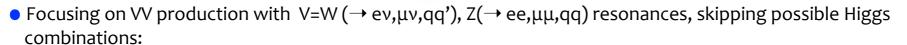


.. no excess observed by ATLAS in other final states or by CMS in all final states.





# Di-Boson F excess! s



Sequential Standard Model (SSM), Extended Gauge Model (EGM) W' → WZ decays,

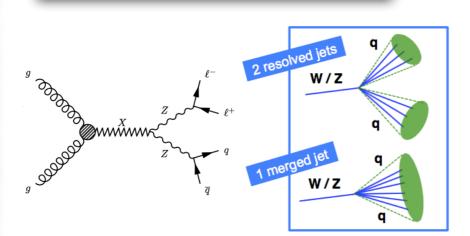
• RS graviton  $G \rightarrow WW$ ,  $G \rightarrow ZZ$  decay.

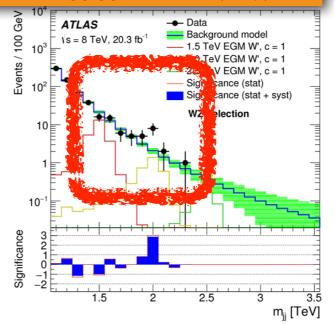
Again, boosted topologies for V→qq

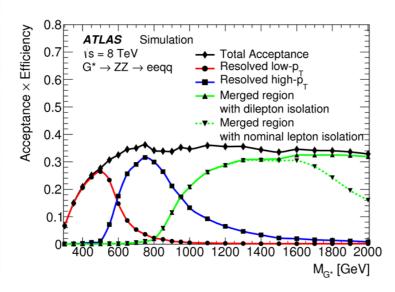
at high mass resonances

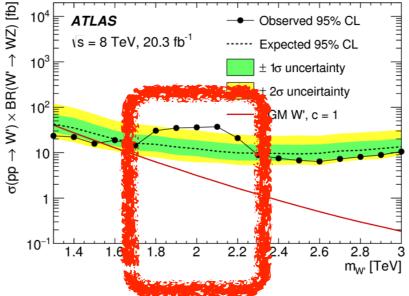
... an (3.4  $\sigma$  local, 2.5  $\sigma$  global) excess observed at M ~ 2 TeV by ATLAS in the

 $VV \rightarrow qq qq final state (only)$ 

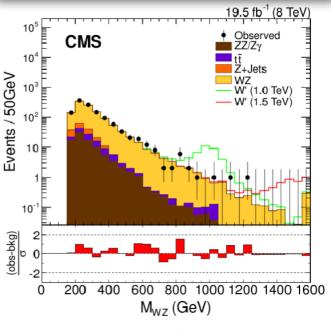


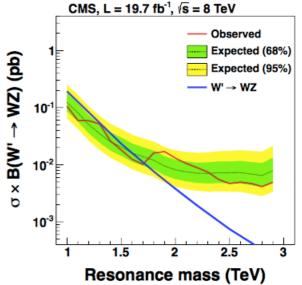






.. no excess observed by ATLAS in other final states or by CMS in all final states.

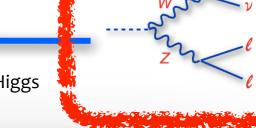




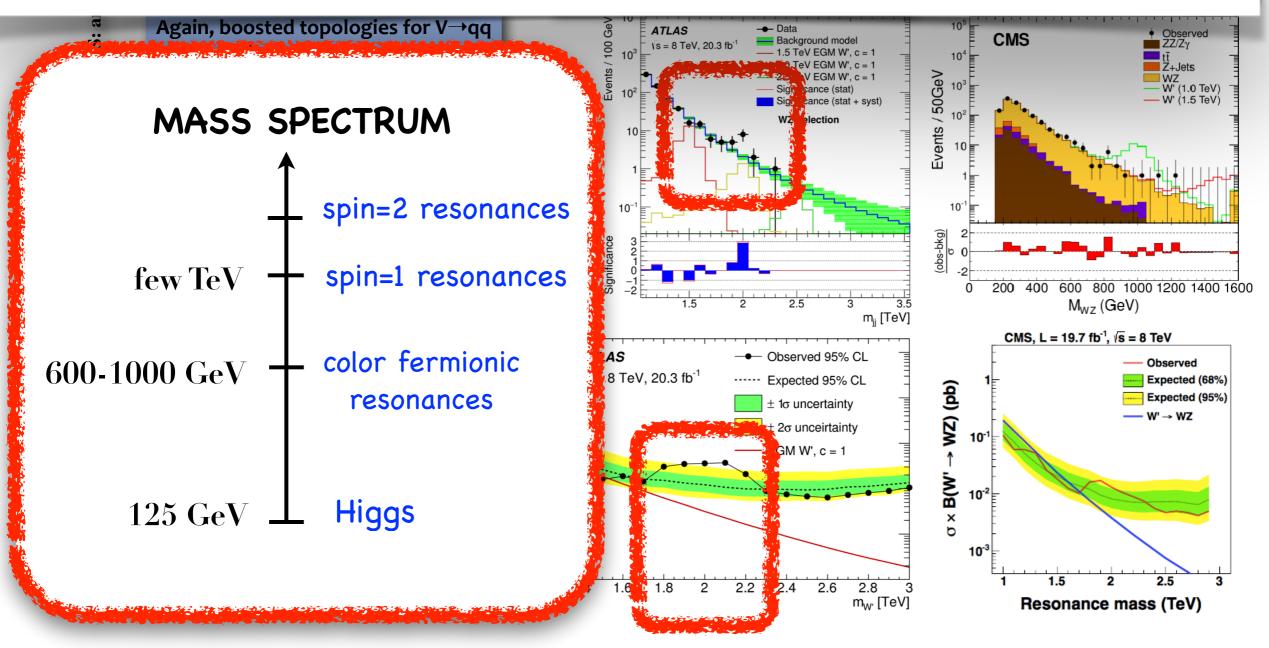
arXiv.1405 3447

## Di-Boson F excess? s

- Focusing on VV production with V=W ( $\rightarrow$  ev, $\mu\nu$ ,qq'), Z( $\rightarrow$  ee, $\mu\mu$ ,qq) resonances, skipping possible Higgs combinations:



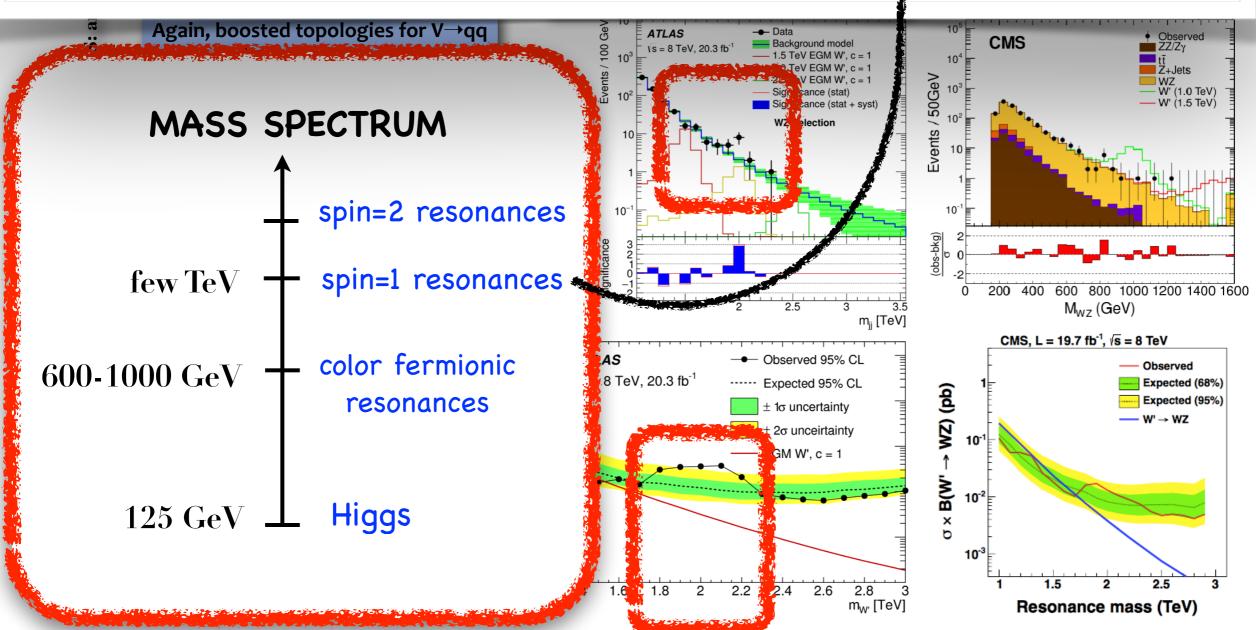
## Expected in composite Higgs (or warped extra-dim) !



## Di-Boson F excess? s

- Focusing on VV production with V=W ( $\rightarrow$  ev, $\mu\nu$ ,qq'), Z( $\rightarrow$  ee, $\mu\mu$ ,qq) resonances, skipping rossible Higgs combinations:

Expected in composite Higgs (or warped extra-dim) !



## New-Physics at the TeV

Pros Cons Cons

Hierarchy problem

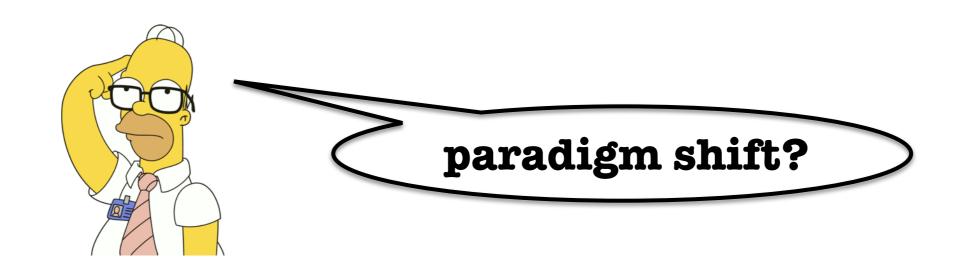
No new particles seen,
no new flavor-violations seen,
no deviations on Higgs couplings seen,
no deviations on Z/W couplings seen,
no WIMP detected,
no EDMs seen,

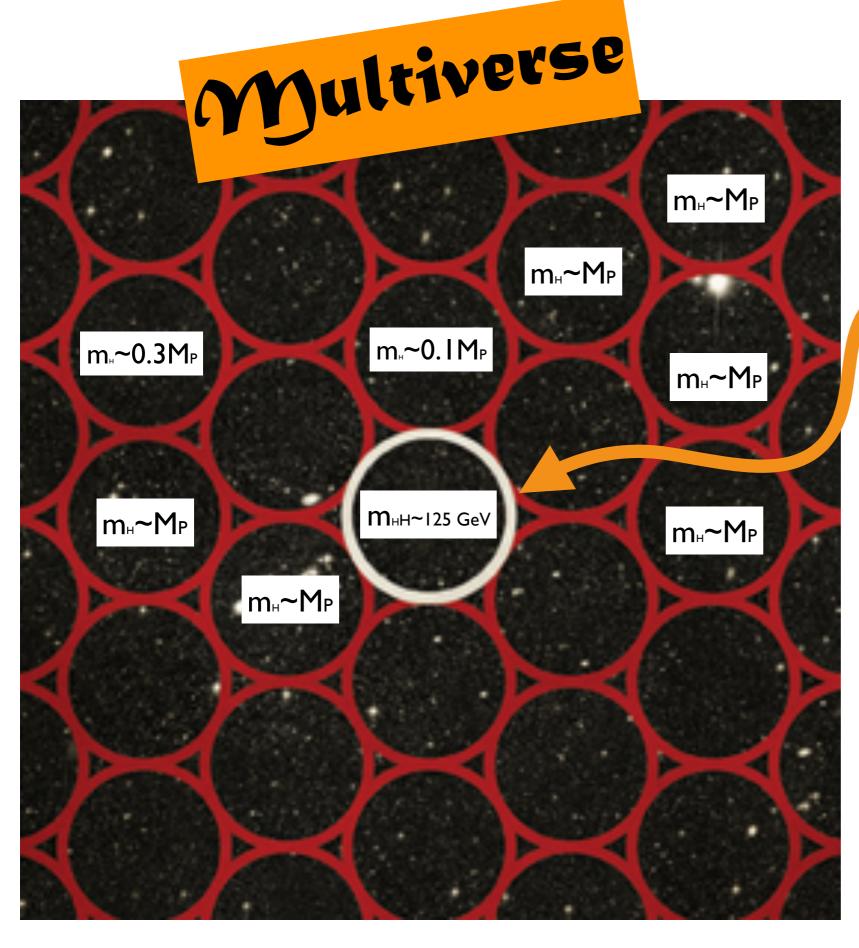
## New-Physics at the TeV

**Pros** Cons

Hierarchy problem

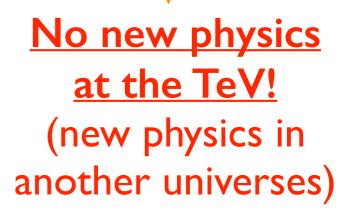
No new particles seen,
no new flavor-violations seen,
no deviations on Higgs couplings seen,
no deviations on Z/W couplings seen,
no WIMP detected,
no EDMs seen,





Our Universe is
very delicate:
Change the SM parameters
and could be uninhabitable

"Natural", since only we can "live" in a Universe with these "fine-tuned" parameters



## Another new Idea for the Hierarchy Problem:

"Relaxation" mechanism

P.W. Graham, D.E. Kaplan, S.Rajendran arXiv:1504.07551

(see also earlier work by Abbott 85, G.Dvali, A.Vilenkin 04, G.Dvali 06)

Higgs-mass parameter — Field-dependent Higgs mass

$$m_H^2 |H|^2$$
  $m_H^2 (\phi) |H|^2$ 

## Another new Idea for the Hierarchy Problem:

#### "Relaxation" mechanism

P.W. Graham, D.E. Kaplan, S.Rajendran arXiv:1504.07551

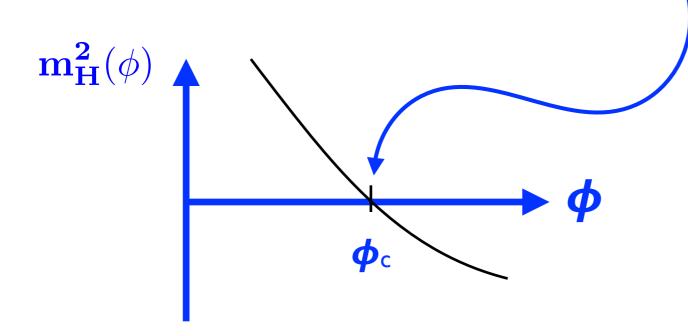
(see also earlier work by Abbott 85, G.Dvali, A.Vilenkin 04, G.Dvali 06)

#### Higgs-mass parameter

#### Field-dependent Higgs mass

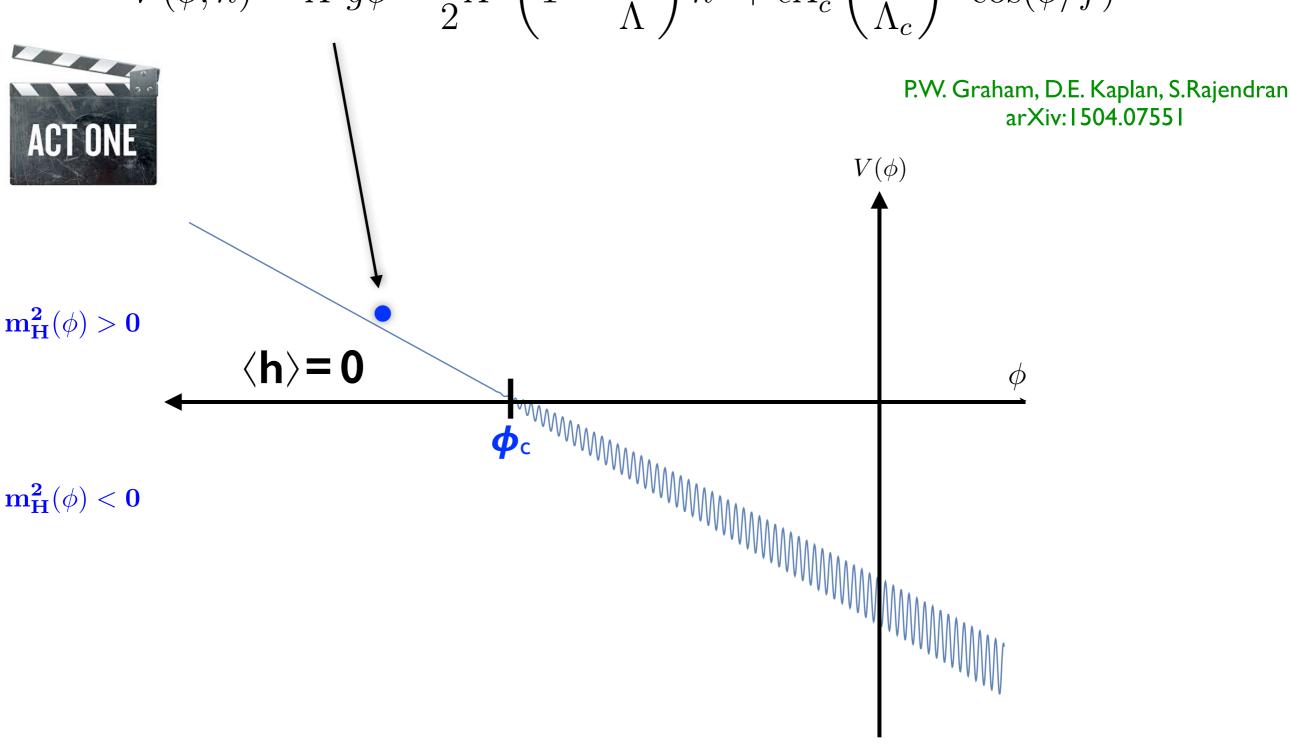
$$m_H^2|H|^2 \qquad \qquad m_H^2(\phi)|H|^2$$

 ${\pmb \phi}$  must get a value where  $m_H^2(\phi) \ll M_P^2$ 

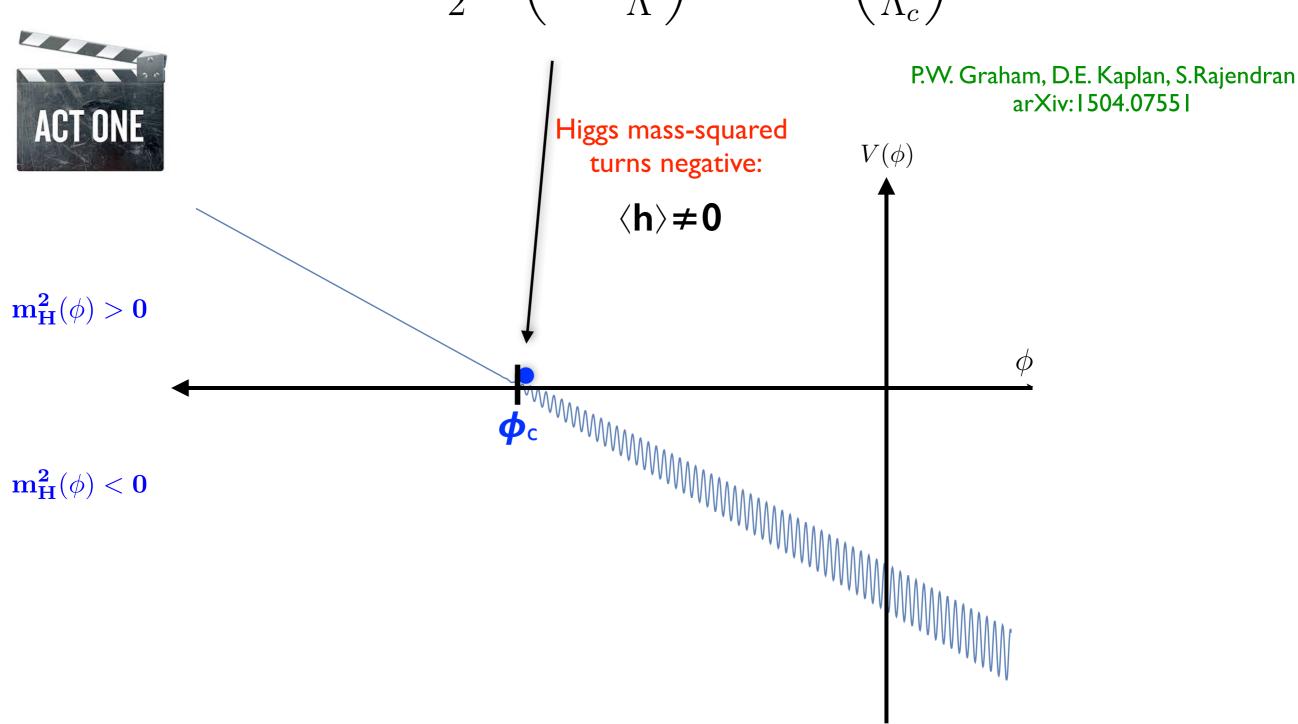


it must arise from a "clever" dynamical interplay between H and  $\phi$ 

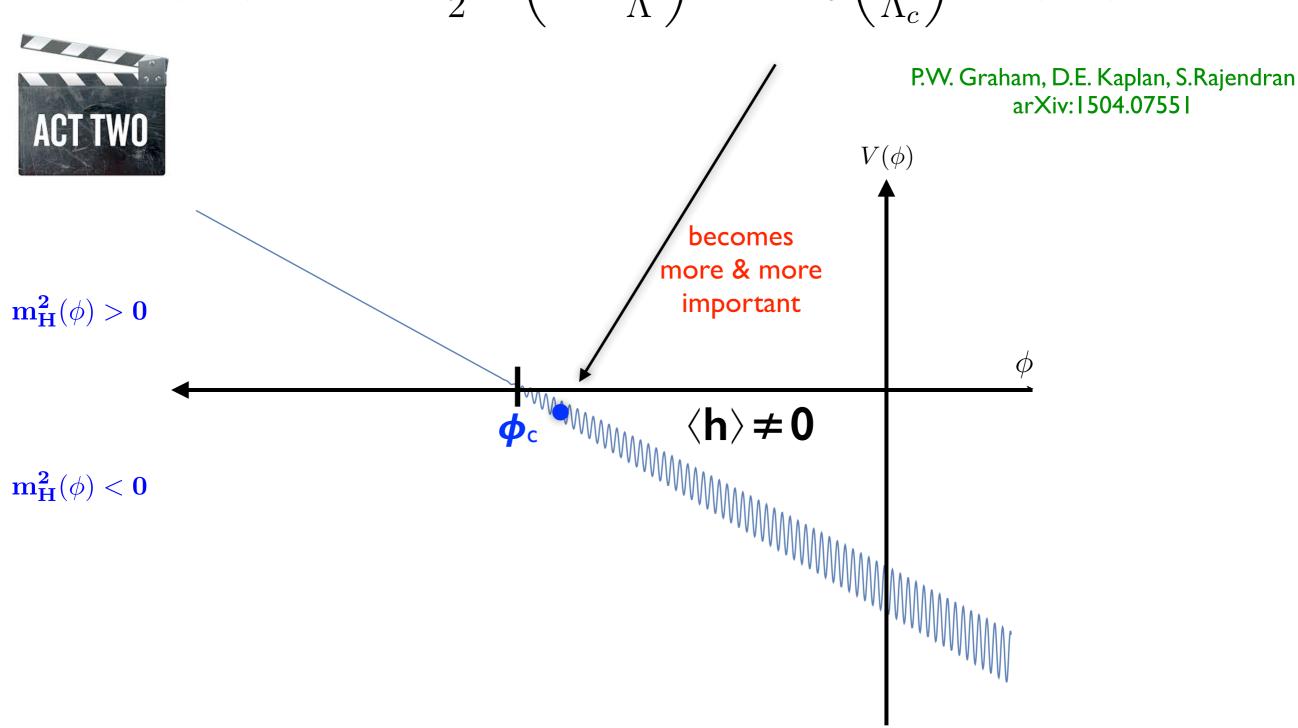
$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g \phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f)$$



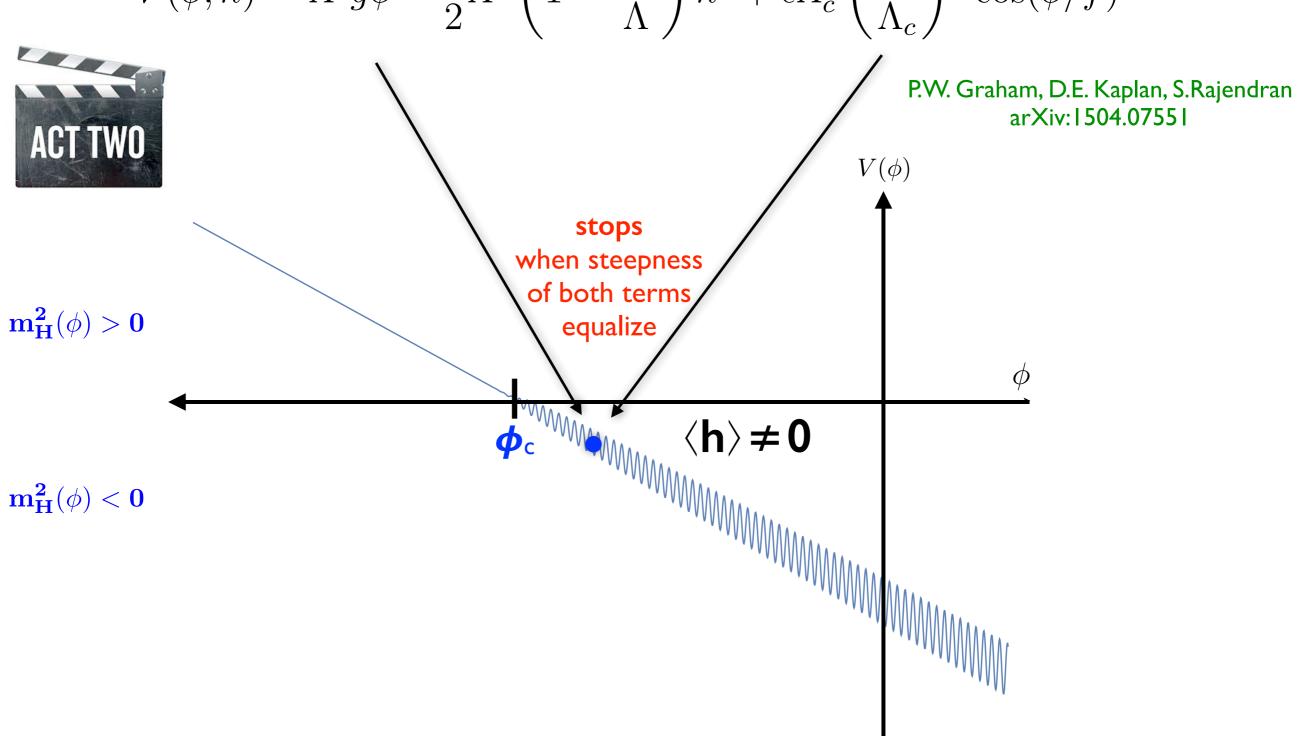
$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g \phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f)$$



$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g \phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f)$$



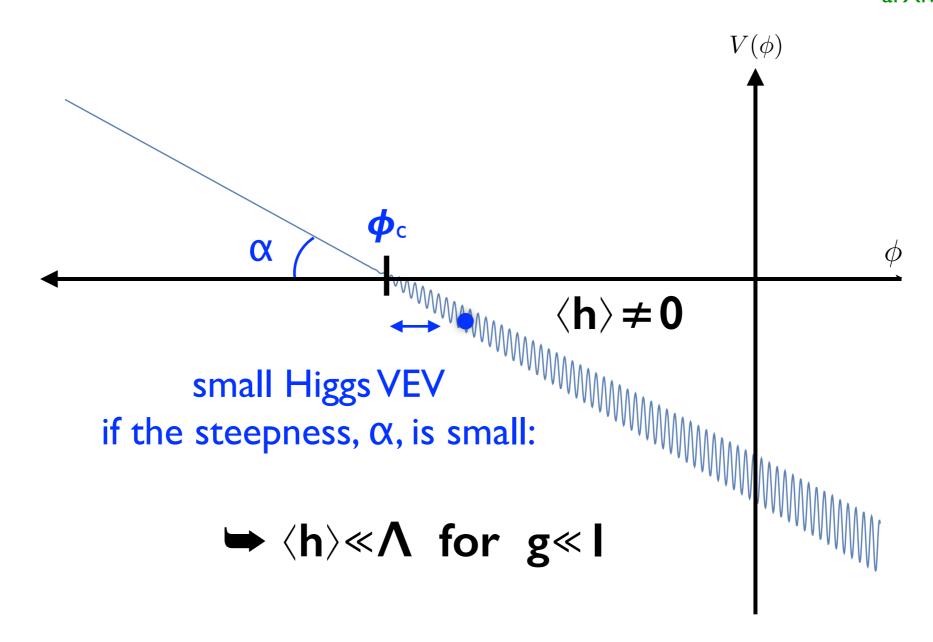
$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g \phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f)$$



#### Cosmological evolution:

$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g \phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f)$$

P.W. Graham, D.E. Kaplan, S.Rajendran arXiv:1504.07551



## Higgs (h) & axion-like ( $\phi$ ) interplay:

$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g \phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f)$$

**\Lambda**: cutoff of the theory

 $\Lambda_c$ : scale that originates the periodic term

#### **Spurions:**

 $g \ll I$ : breaking shift symmetry  $\phi \rightarrow \phi + c$ 

 $\epsilon \ll I$ : breaking of shift symmetry, respecting  $\phi \rightarrow \phi + 2\pi f$ ,  $\phi \rightarrow -\phi$ 

potential stable under radiative corrections!

## Tuning the initial conditions?



#### Tuning the initial conditions?



No, if slow rolling due to a friction: possible in the inflationary epoch! (Hubble friction)

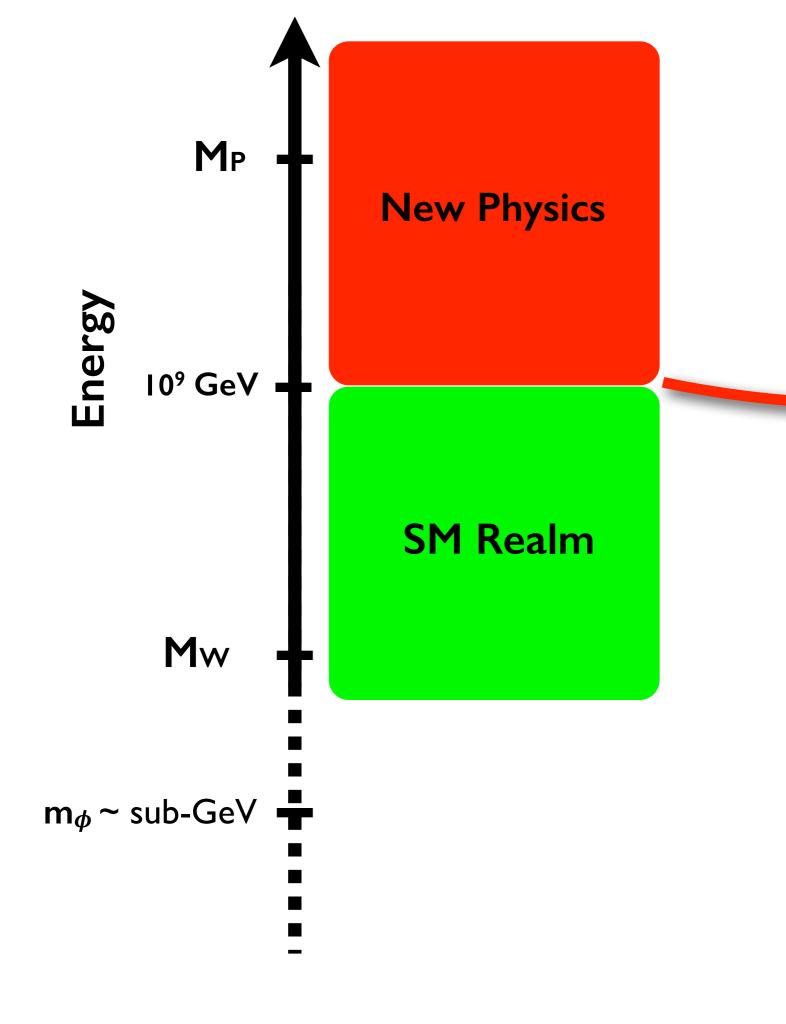
#### Tuning the initial conditions?



No, if slow rolling due to a friction: possible in the inflationary epoch! (Hubble friction)

Long period of inflation needed, in order for  $\phi$  to "scan" large ranges of the Higgs mass

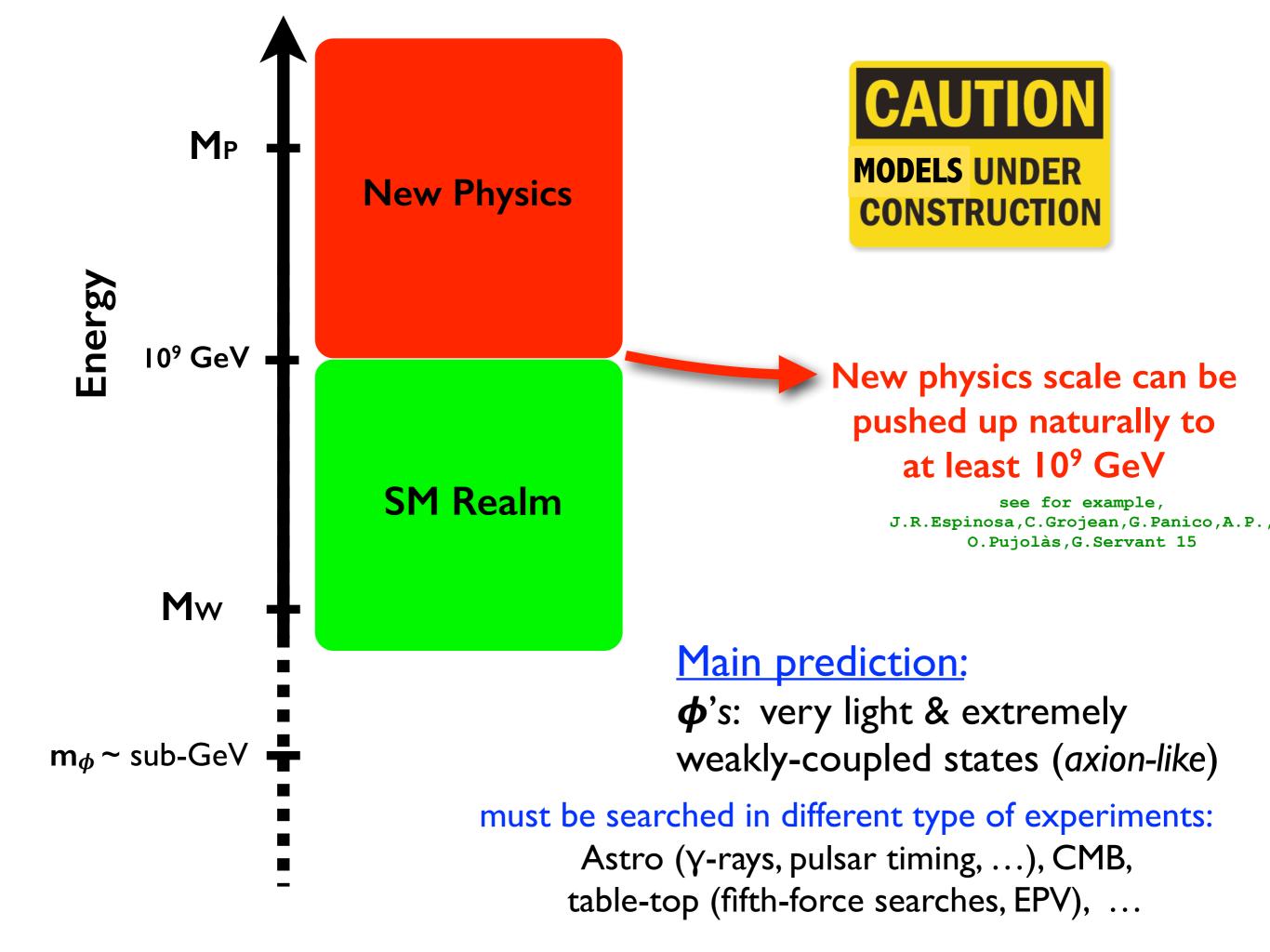
e-folds needed:  $N_e \gtrsim \frac{H_I^2}{g^2\Lambda^2} \sim 10^{40}$ 





New physics scale can be pushed up naturally to at least 109 GeV

see for example,
J.R.Espinosa, C.Grojean, G.Panico, A.P.
O.Pujolàs, G.Servant 15



## Conclusions

◆ After LHC run I → the SM has been completed

No need for anything else
 (at least) up to around the Planck scale

End of no-lose theorems for discovery at the TeV



## Conclusions

After LHC run I is the SM has been completed

No need for anything else (at least) up to around the Planck scale



End of no-lose theorems for discovery at the TeV

- We start a very different phase in particle physics:
   BSM only motivated by the unnaturalness of the SM!
- LHC run 2: Important jump in the energy
  - right chances for new discoveries.

Main one: Learn about the origin of Higgs potential (EWSB origin)

Either supersymmetry, new strong-dynamics or something else

A new era begins...

## Conclusions

After LHC run I is the SM has been completed

No need for anything else (at least) up to around the Planck scale



End of no-lose theorems for discovery at the TeV

- We start a very different phase in particle physics:
   BSM only motivated by the unnaturalness of the SM!
- LHC run 2: Important jump in the energy
  - right chances for new discoveries.

Main one: Learn about the origin of Higgs potential (EWSB origin)

Either supersymmetry, new strong-dynamics or something else

#### A new era begins...

Epilogue: Don't be afraid of null results: As Michelson-Morley experiment, also null results (from well-motivated experiments) can lead to a change of paradigm